

Supersymmetric cascade decays

based on: [B. K. Gjelsten](#), D. J. Miller, P. Osland

[JHEP 12 \(2004\) 003 \[hep-ph/0410303\]](#), [hep-ph/0501033](#)

Per Osland, University of Bergen

Outline

- Introduction
- Graphic overview of parameter space
- Kinematical endpoints (squark chain)
- Inverting endpoint formulas
- Fitting MC data: precision
- 10,000 LHC experiments: precision
- LC input: high precision
- gluino chain (b-tagging)

Supersymmetry may be realized at the LHC

Unstable particles produced copiously,
cascade decays, e.g.

$$\tilde{q} \rightarrow \tilde{\chi}_2^0 q \rightarrow \tilde{l} l q \rightarrow \tilde{\chi}_1^0 l l q$$

Challenge: determine masses with high precision

Refs: Baer et al, hep-ph/9512383; Hinchliffe et al, hep-ph/9610544;
Bachacou et al, hep-ph/9907518; Polesello, ATLAS Int Note 1997;
Allanach et al, hep-ph/0007009; Gjelsten et al, ATLAS Note 2004;
Chiorboli, Tricomi, CMS Note 2004

mSUGRA (CMSSM)

- Unification at high energies, fewer parameters

$$m_0 \quad m_{1/2} \quad A_0 \quad \tan \beta \quad \text{sign } \mu$$

- Snowmass Points and Slopes: [Allanach et al, hep-ph/0202233](#): SPS 1a, SPS 1b, SPS 3, SPS 5.

Other set: Battaglia et al, hep-ph/0306219

- WMAP constraints: Bennett et al, astro-ph/0302207;
Spergel et al, astro-ph/0302209

footnote: Benchmark Points

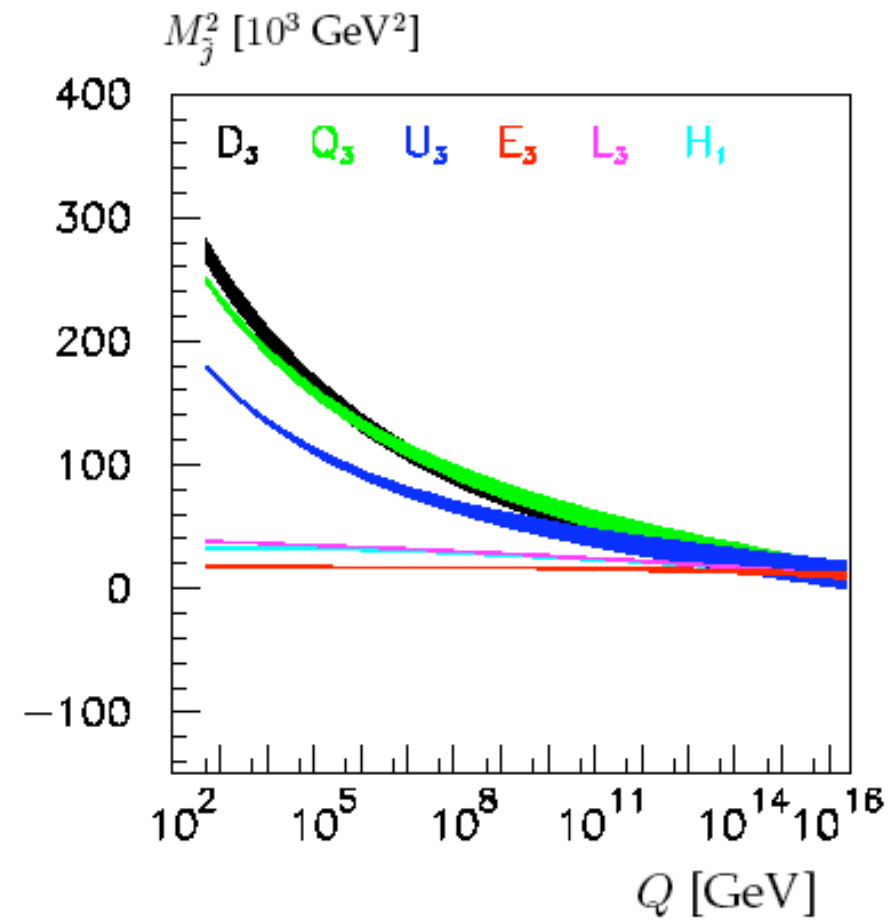
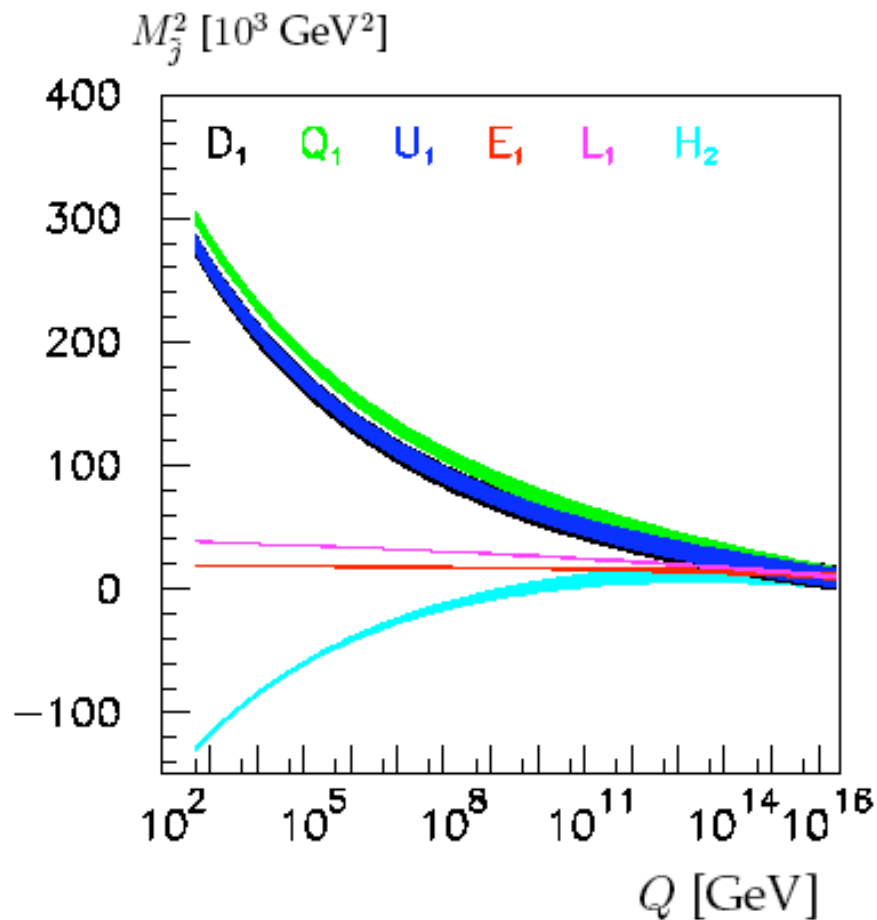
- LHC Points ('SUGRA'): [Hinchliffe et al, hep-ph/9610544](#): [Point 1](#), [Point 2](#), [Point 3](#), [Point 4](#), [Point 5](#)
- Post-LEP Benchmarks ('CMSSM'): [Battaglia et al, hep-ph/0106204](#): [A](#), [B](#), [C](#), ..., [M](#)
- ➔ ● Snowmass Points and Slopes ('mSUGRA'): [Allanach et al, hep-ph/0202233](#): [SPS 1a](#), [SPS 1b](#), [SPS 2](#), [SPS 3](#), [SPS 4](#), [SPS 5](#), [SPS 6](#), ..., [SPS 9](#)
- Post-WMAP Benchmarks ('CMSSM'): [Ellis et al, hep-ph/0303043](#), [Battaglia et al hep-ph/0306219](#): [A'](#), [B'](#), [C'](#), ..., [M'](#)

Mutations: Point 5 \rightarrow B \rightarrow SPS 1a \rightarrow B'

Precision in masses allows extrapolation to Unification scale

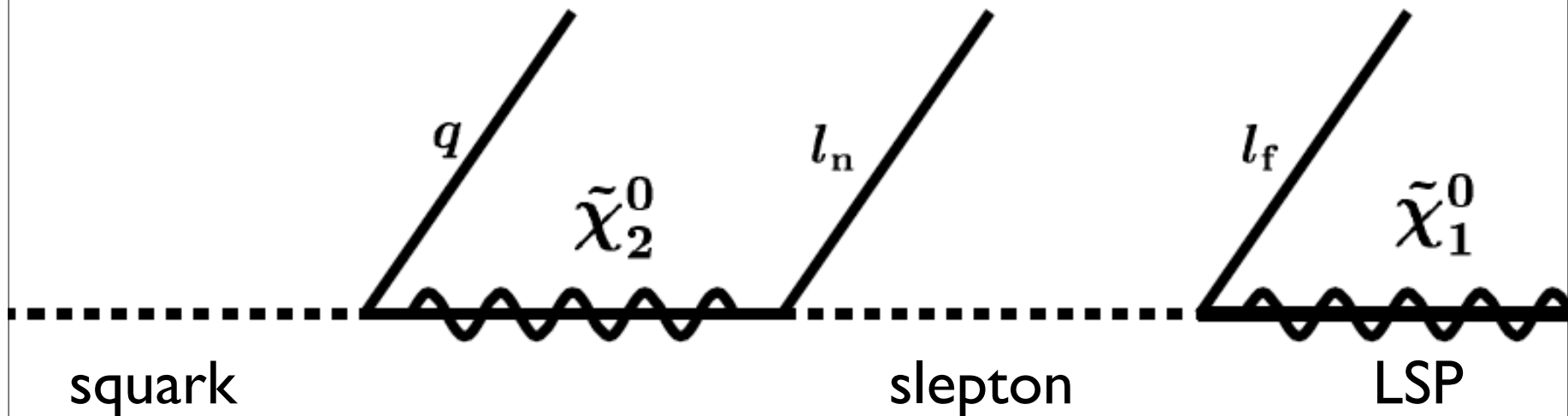
Example: Sfermion mass parameters

Allanach et al, hep-ph/0403133



Precision of order per cent achievable with LHC plus LC

“Easy” SPS Ia squark cascade



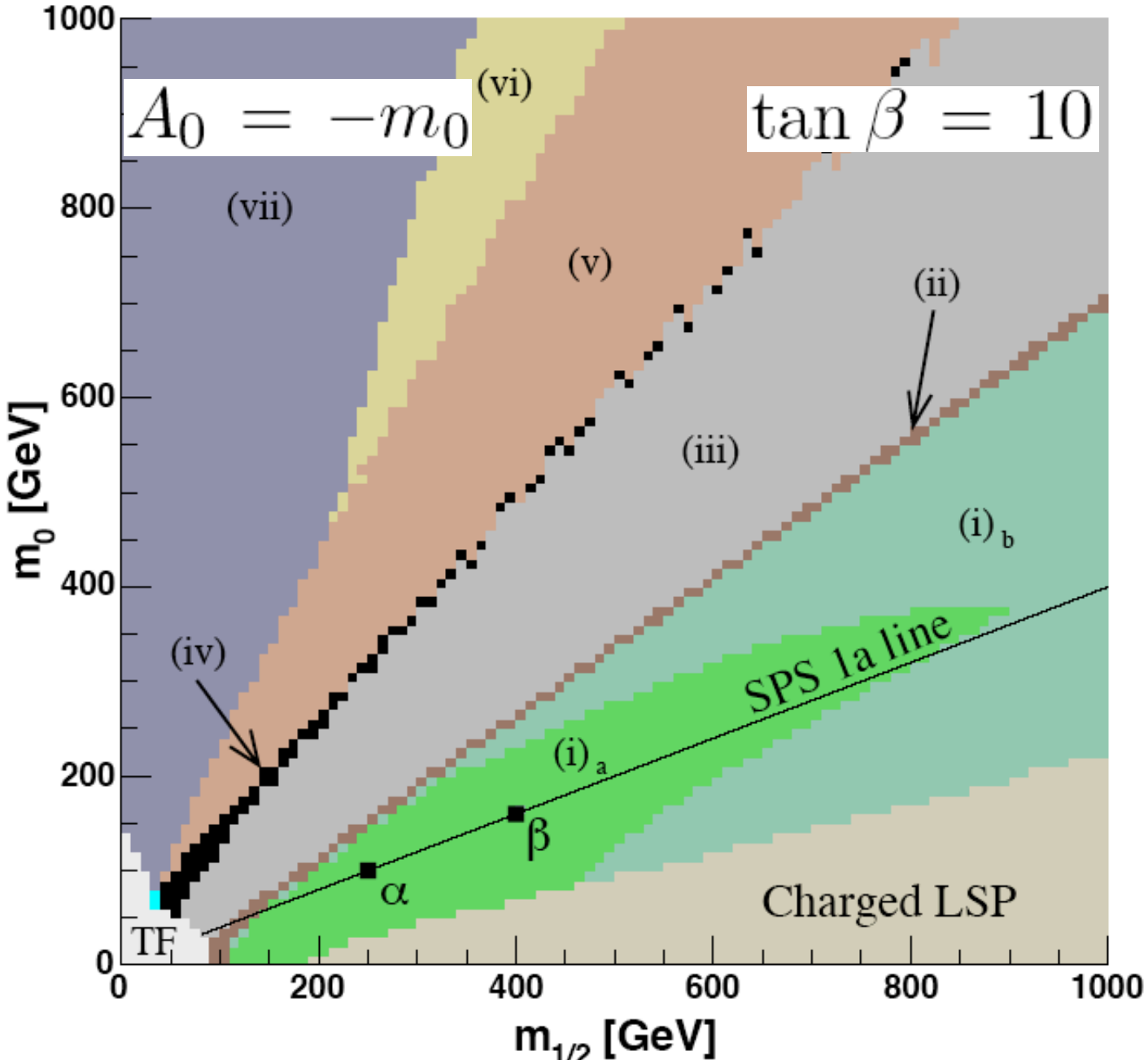
Detect: **quark jet** and **two leptons**

Aim: determine **squark**, **slepton** and **neutralino masses**

Question: Is this mass hierarchy “typical”?

Want “heavy” gluino and “heavy” neutralino $\tilde{\chi}_2^0$

Hierarchies:

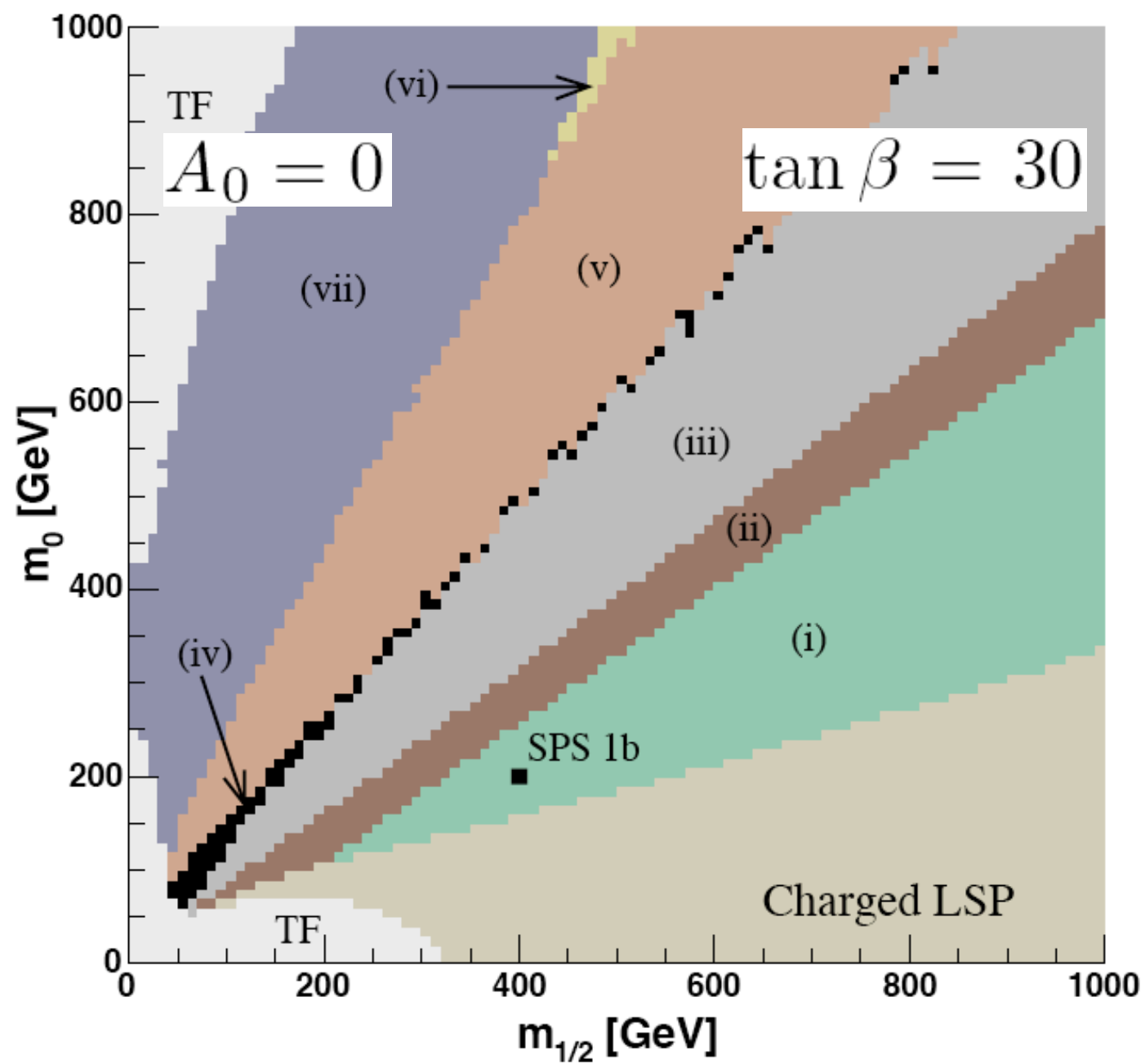


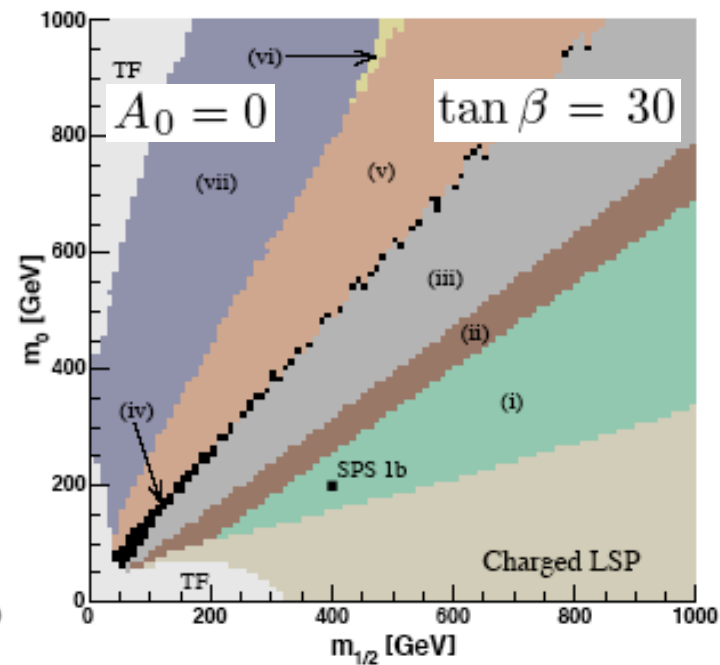
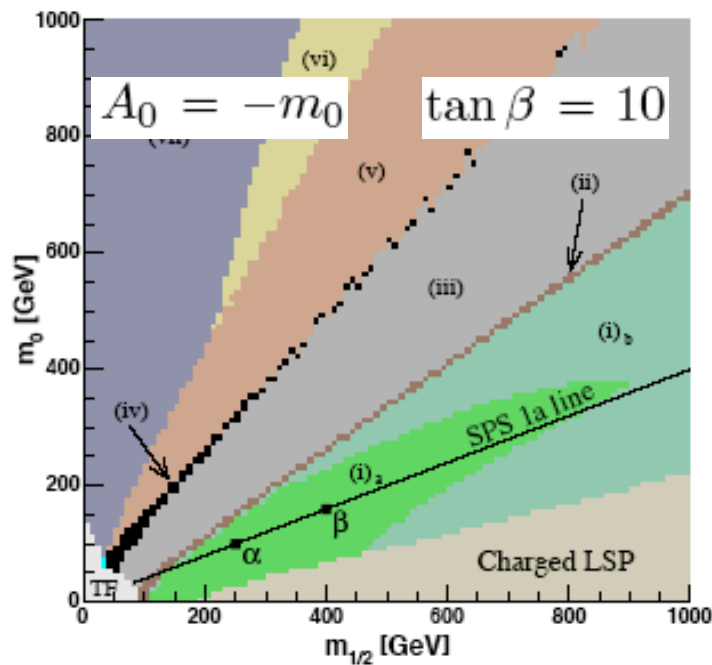
heavy gluino

heavy neutralino

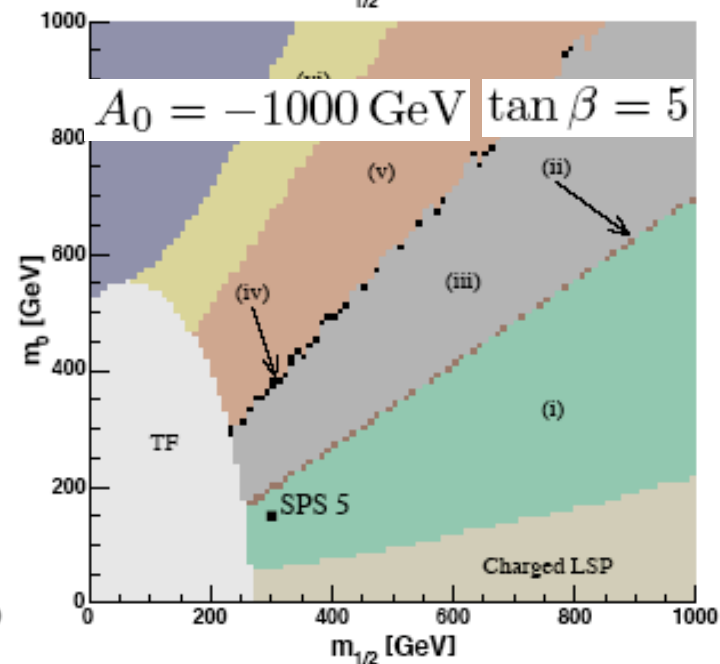
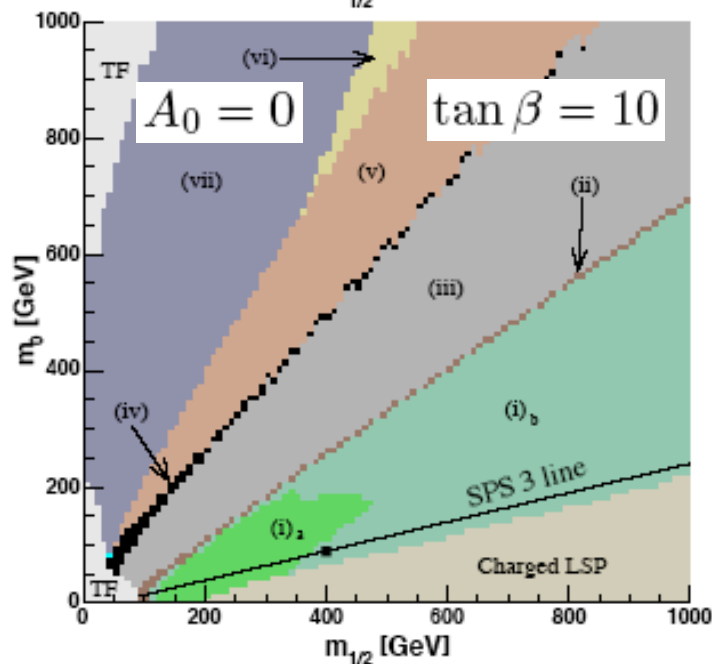
- | | | | |
|-------|---|-----|--|
| (i) | $\tilde{g} > \max(\tilde{d}_L, \tilde{u}_L, \tilde{b}_1, \tilde{t}_1)$ | and | $\tilde{\chi}_2^0 > \max(\tilde{l}_R, \tilde{\tau}_1)$ |
| (ii) | $\tilde{g} > \max(\tilde{d}_L, \tilde{u}_L, \tilde{b}_1, \tilde{t}_1)$ | and | $\tilde{l}_R > \tilde{\chi}_2^0 > \tilde{\tau}_1$ |
| (iii) | $\tilde{g} > \max(\tilde{d}_L, \tilde{u}_L, \tilde{b}_1, \tilde{t}_1)$ | and | $\min(\tilde{l}_R, \tilde{\tau}_1) > \tilde{\chi}_2^0$ |
| (iv) | $\tilde{d}_L > \tilde{g} > \max(\tilde{u}_L, \tilde{b}_1)$ | and | $\min(\tilde{l}_R, \tilde{\tau}_1) > \tilde{\chi}_2^0$ |
| (v) | $\min(\tilde{d}_L, \tilde{u}_L) > \tilde{g} > \tilde{b}_1$ | and | $\min(\tilde{l}_R, \tilde{\tau}_1) > \tilde{\chi}_2^0$ |
| (vi) | $\min(\tilde{d}_L, \tilde{u}_L, \tilde{b}_1) > \tilde{g} > \tilde{t}_1$ | and | $\min(\tilde{l}_R, \tilde{\tau}_1) > \tilde{\chi}_2^0$ |
| (vii) | $\min(\tilde{d}_L, \tilde{u}_L, \tilde{b}_1, \tilde{t}_1) > \tilde{g}$ | and | $\min(\tilde{l}_R, \tilde{\tau}_1) > \tilde{\chi}_2^0$ |

heavy gauginos, $\tilde{g}, \tilde{\chi}_2^0$ lower right





All:
 $\mu > 0$

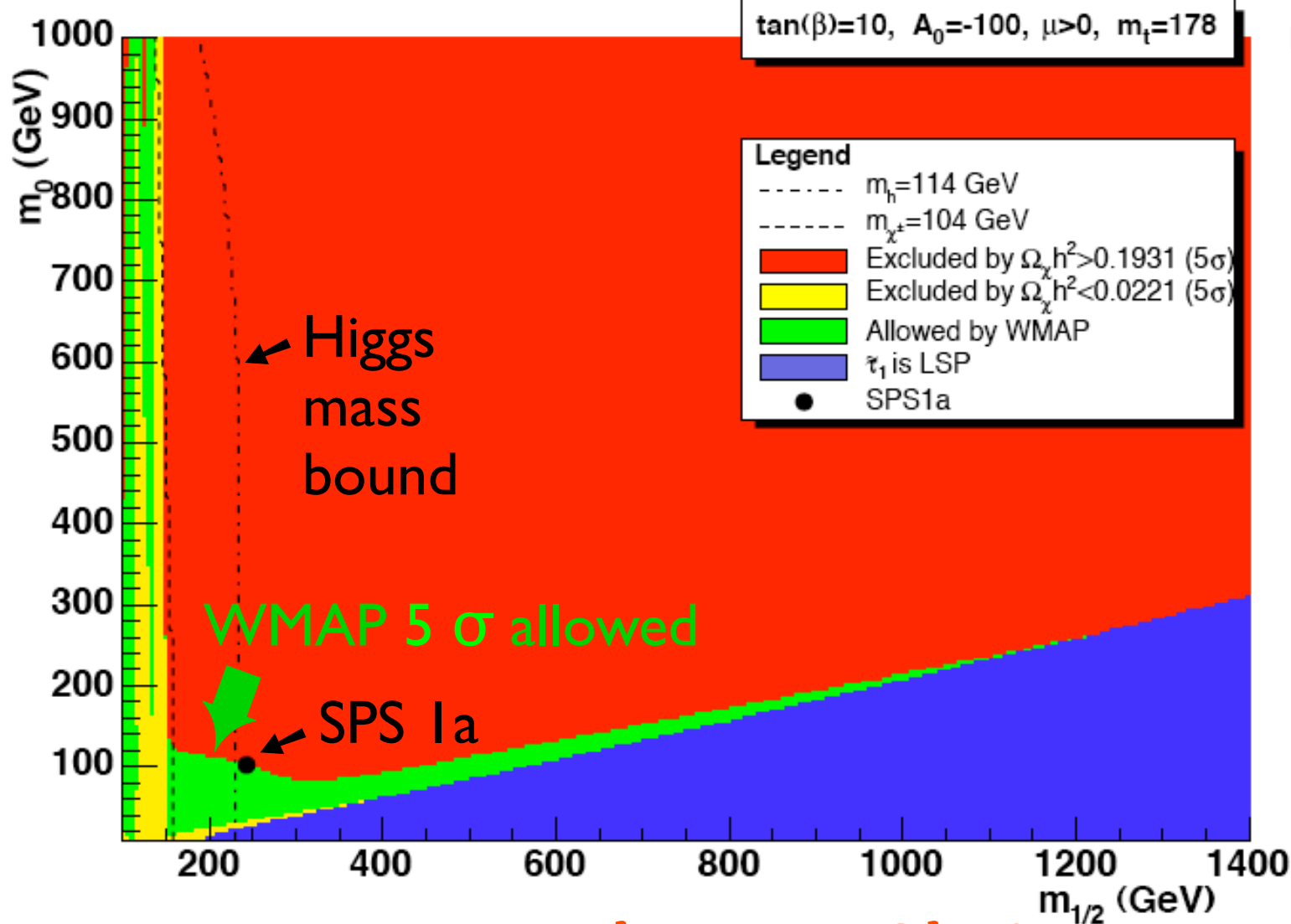


WMAP & LEP constraints

WMAP constraints

from Are Raklev

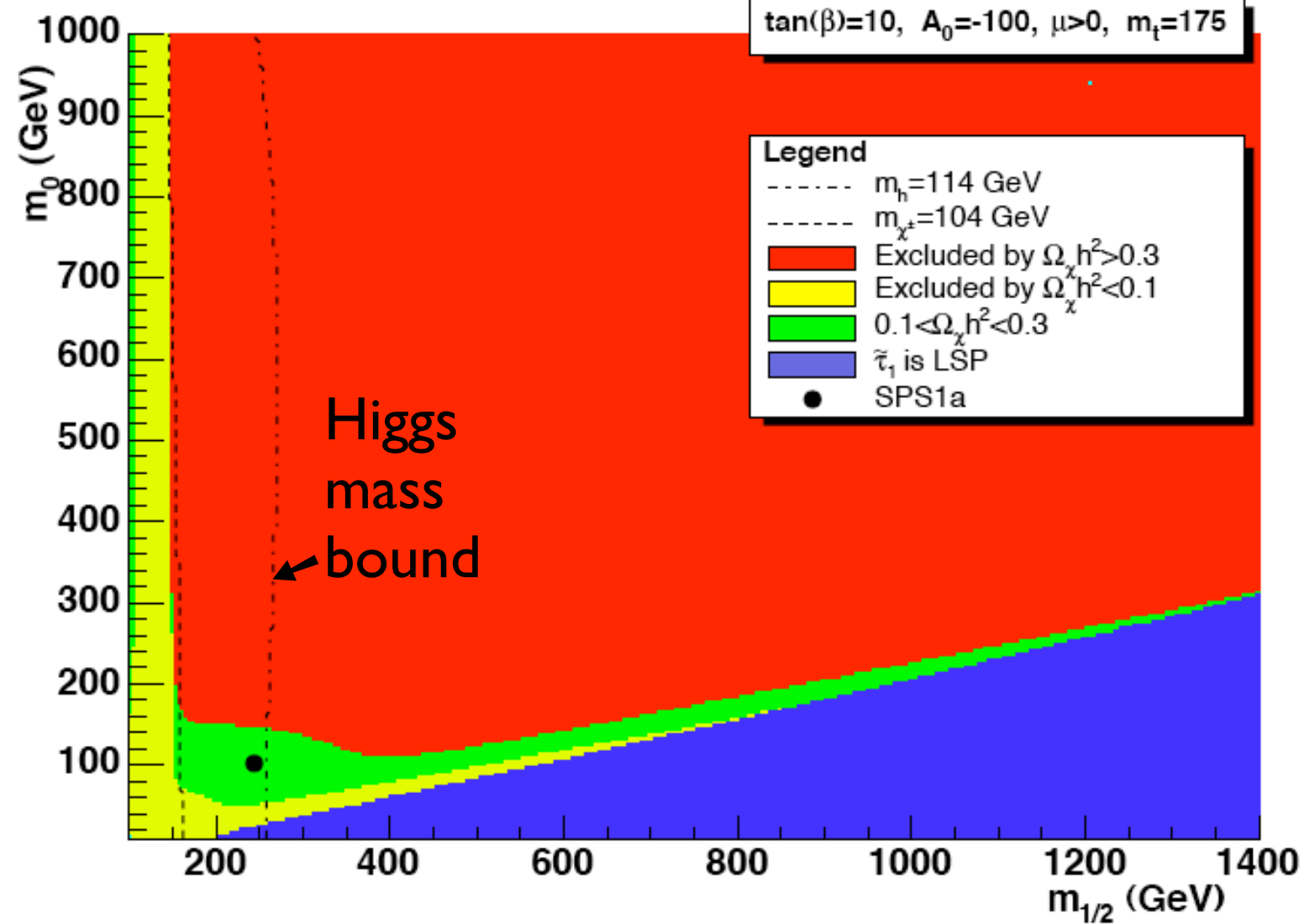
DarkSusy
ISASUSY



Ignore slight conflict!

from Are Raklev

General cosmological constraints



Next question:

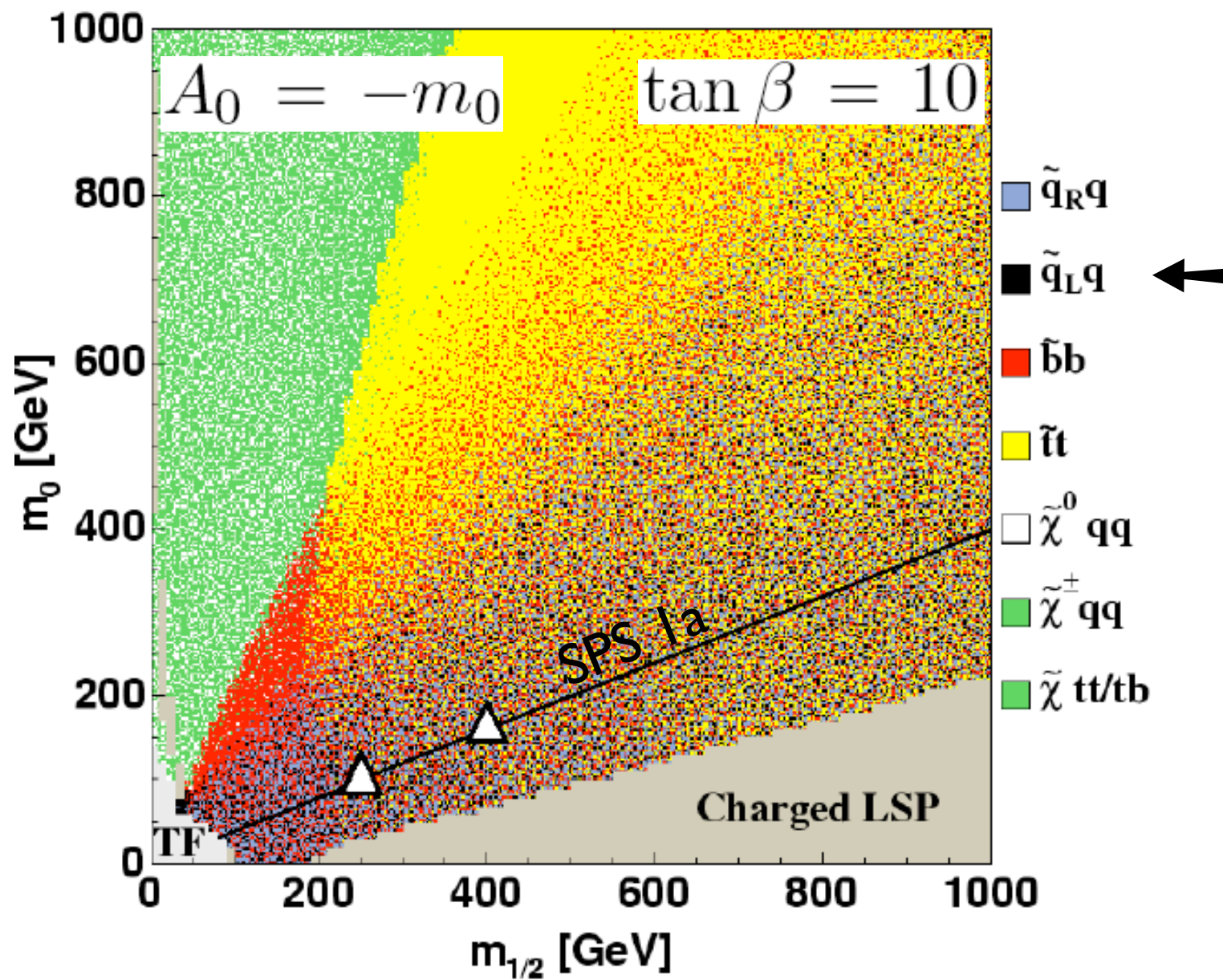
Given “correct” hierarchy,

$$m_{\tilde{g}} > m_{\tilde{q}_L} > m_{\tilde{\chi}_2^0} > m_{\tilde{l}_R} > m_{\tilde{\chi}_1^0}$$

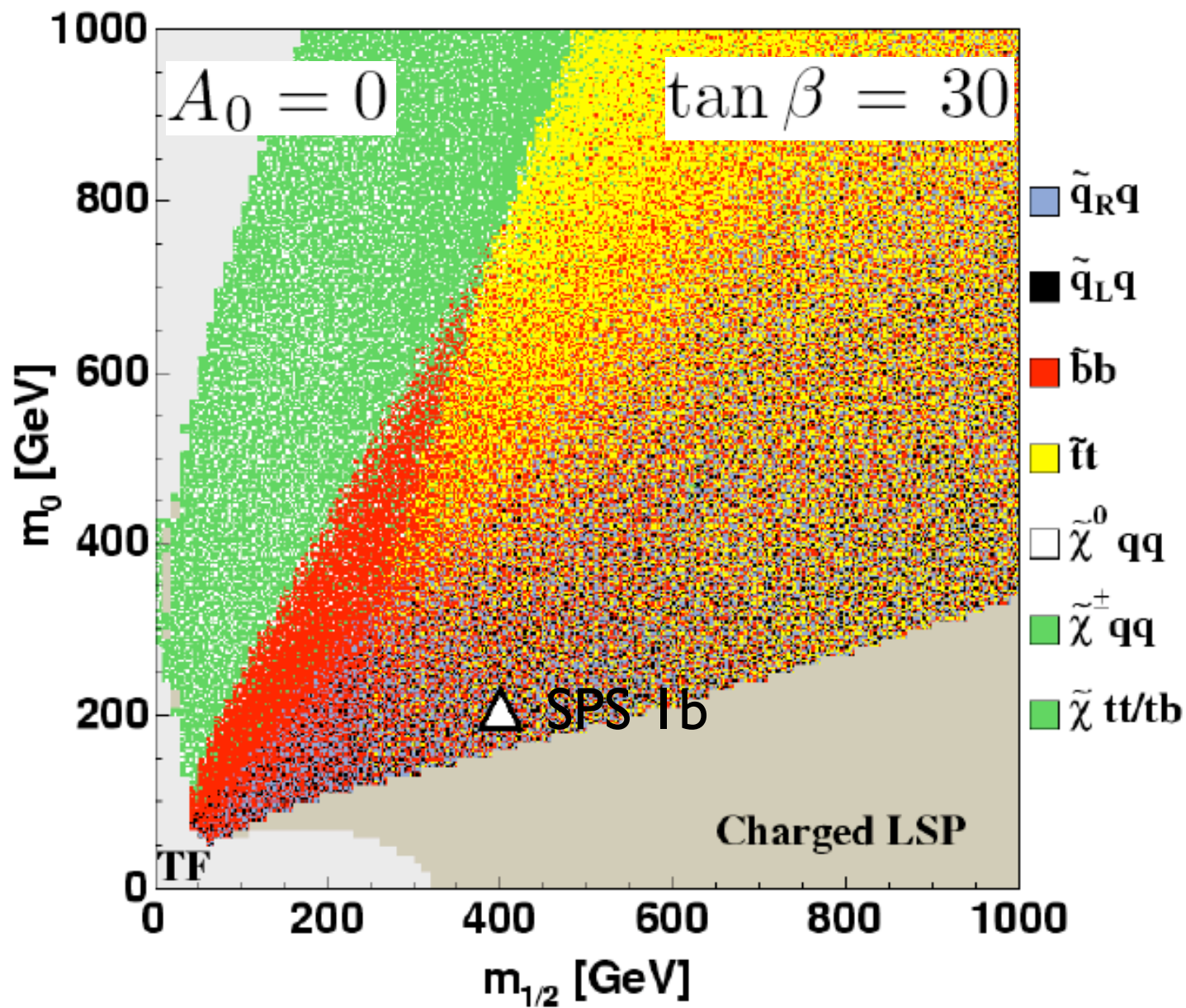
is there enough BR?

- Does the squark have significant BR to neutralino and quark?
- Does the neutralino have significant BR to slepton and lepton?

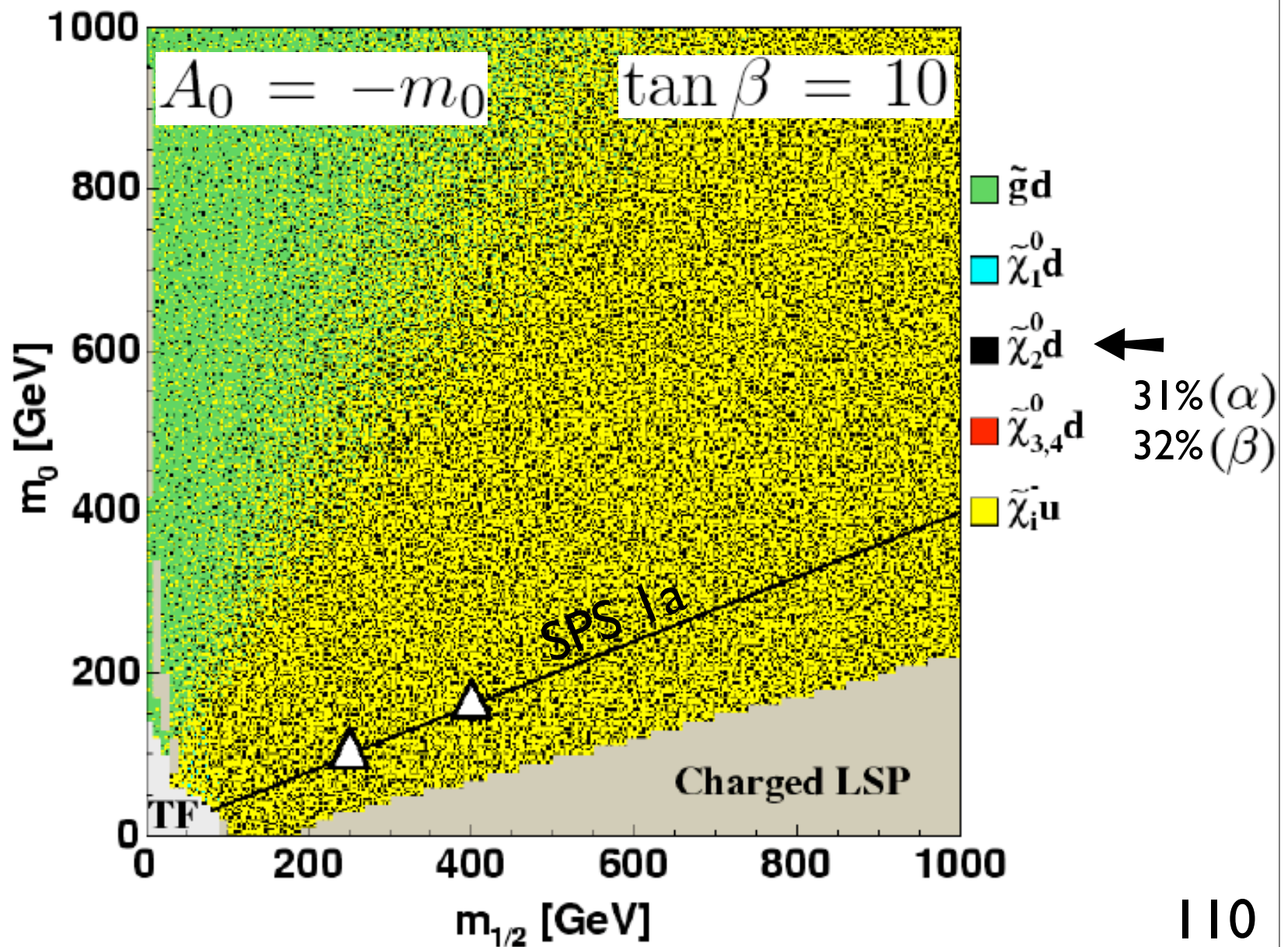
Gluino Branching Ratios (\tilde{g})



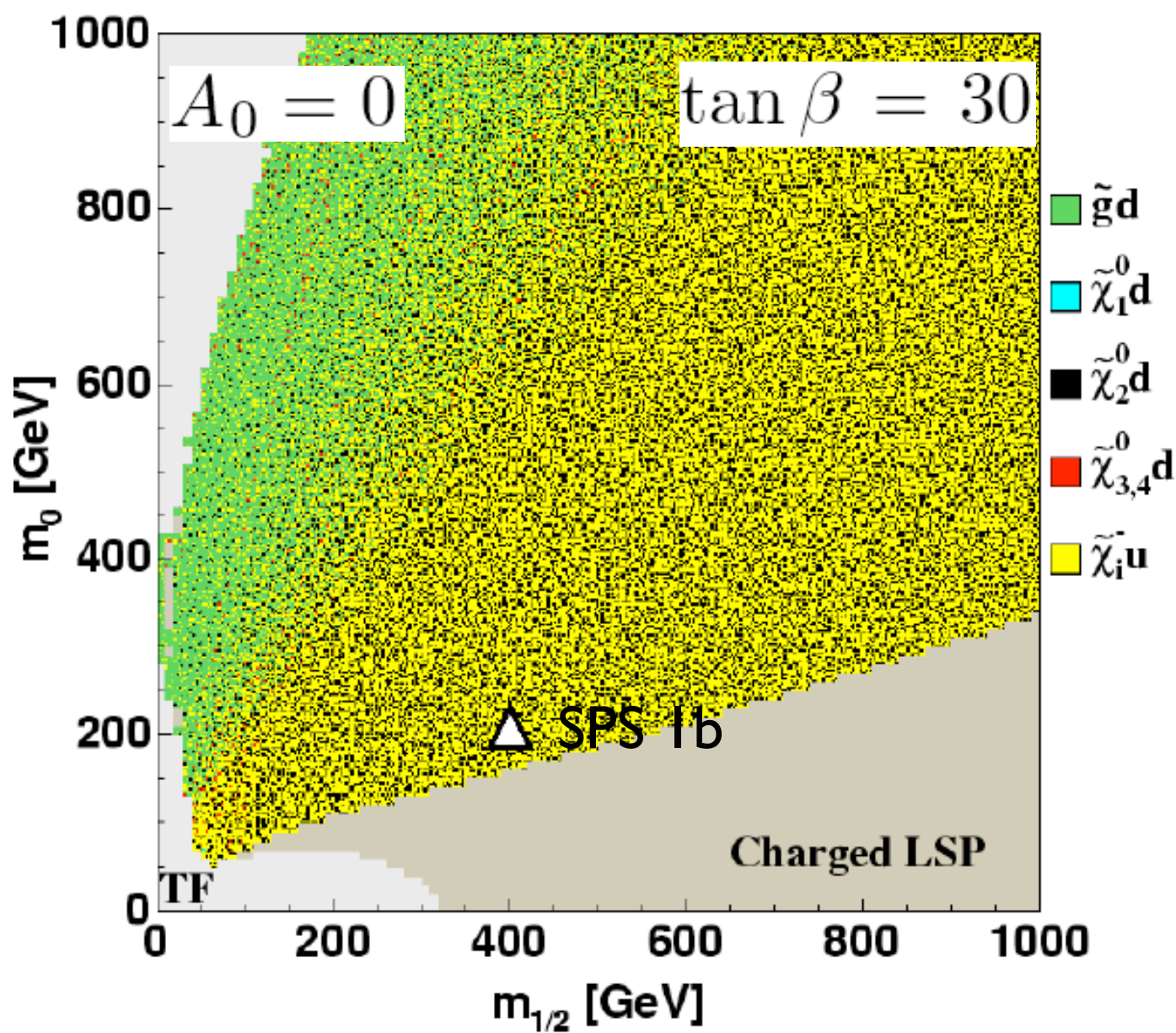
Gluino Branching Ratios (\tilde{g})



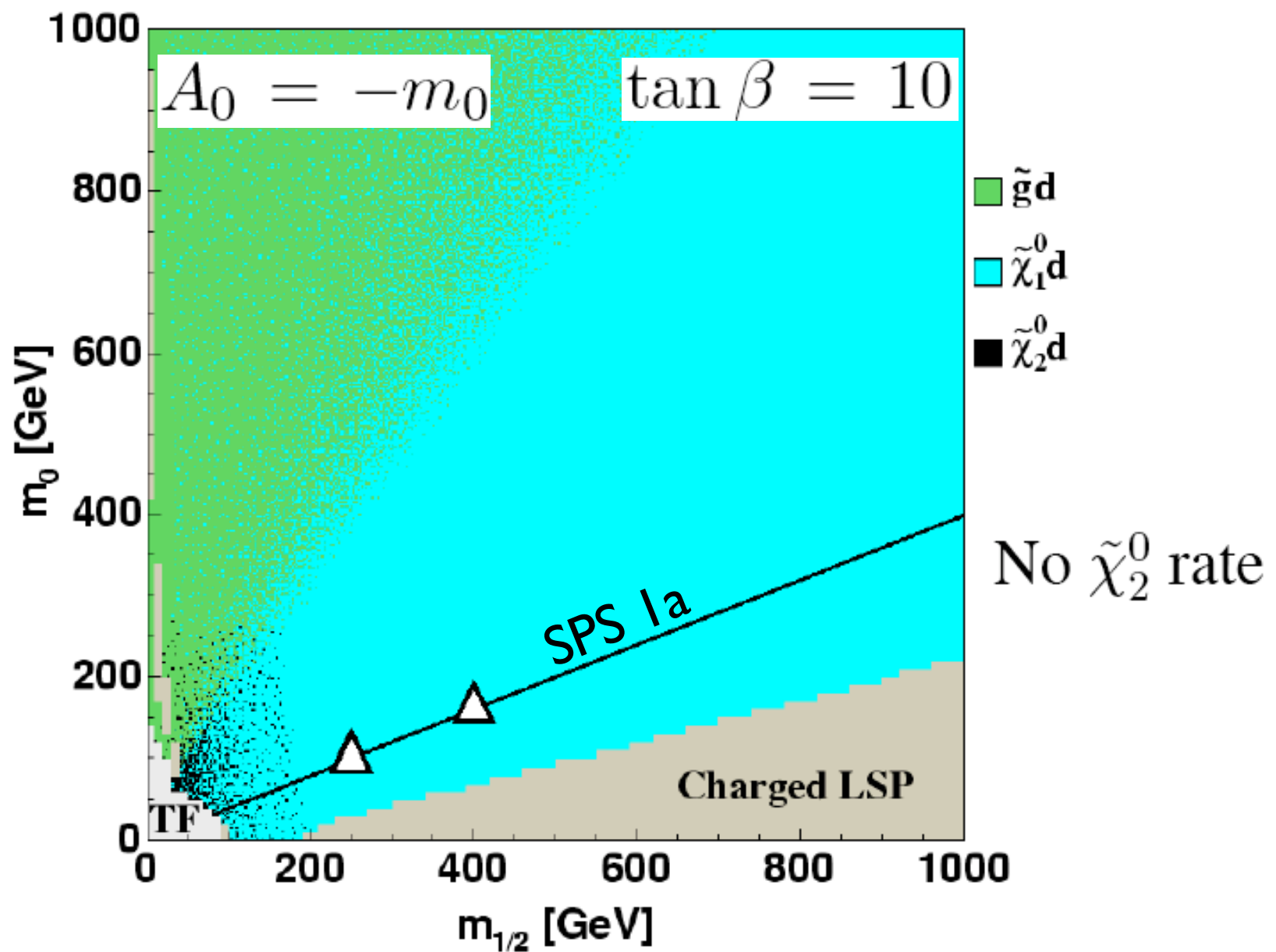
Squark Branching Ratios (\tilde{u}_L)



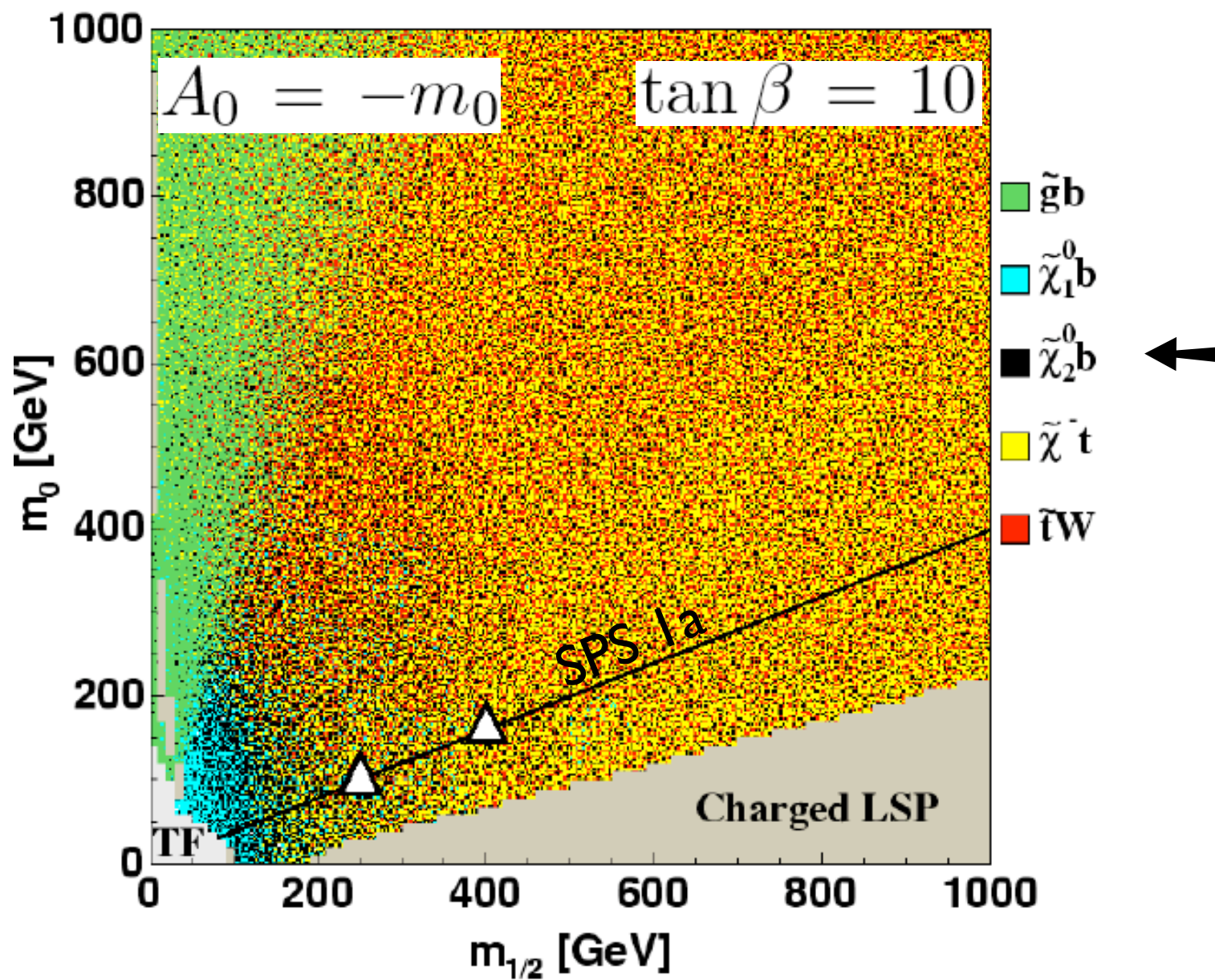
Squark Branching Ratios (\tilde{u}_L)



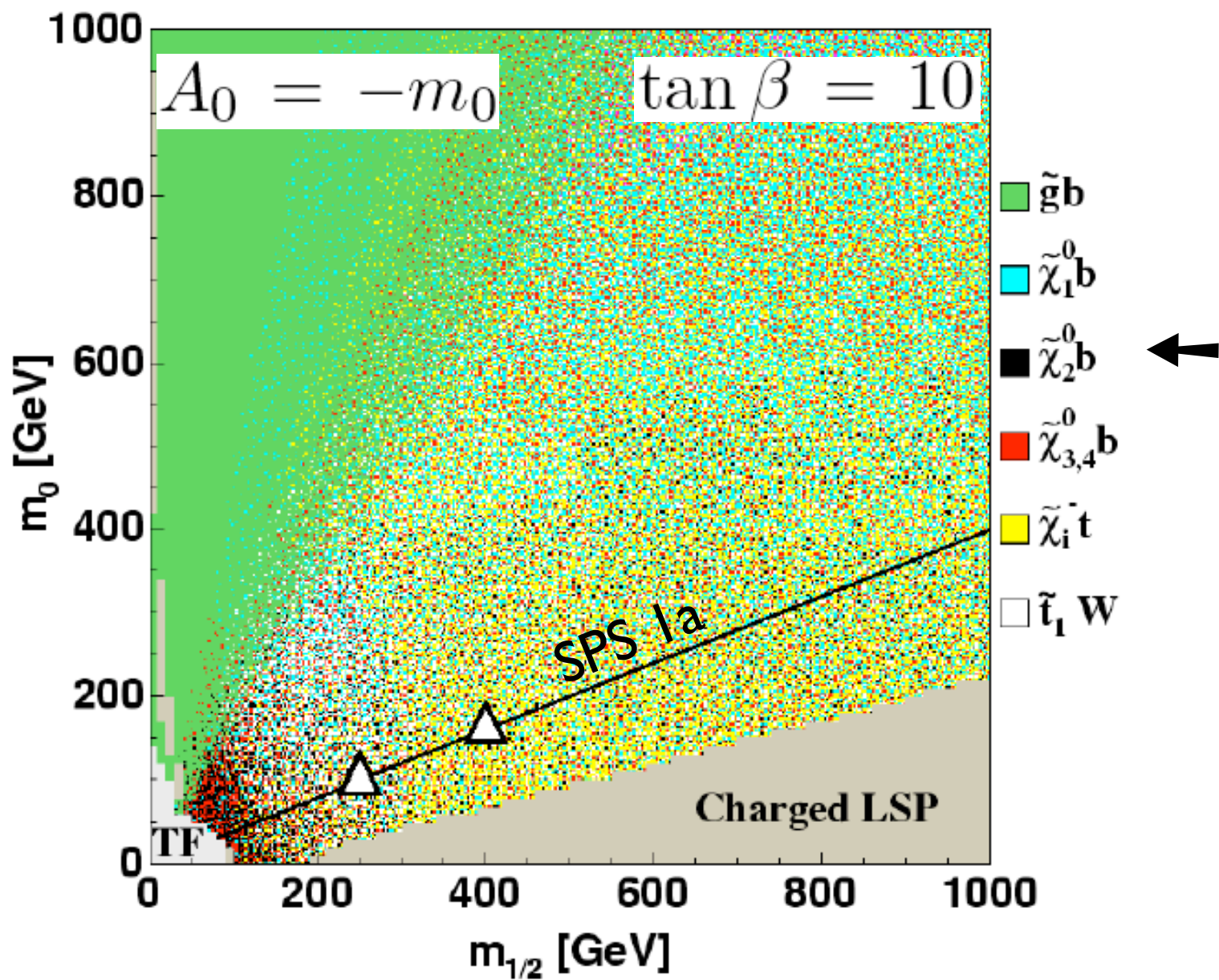
Squark Branching Ratios (~~\tilde{u}_R~~)



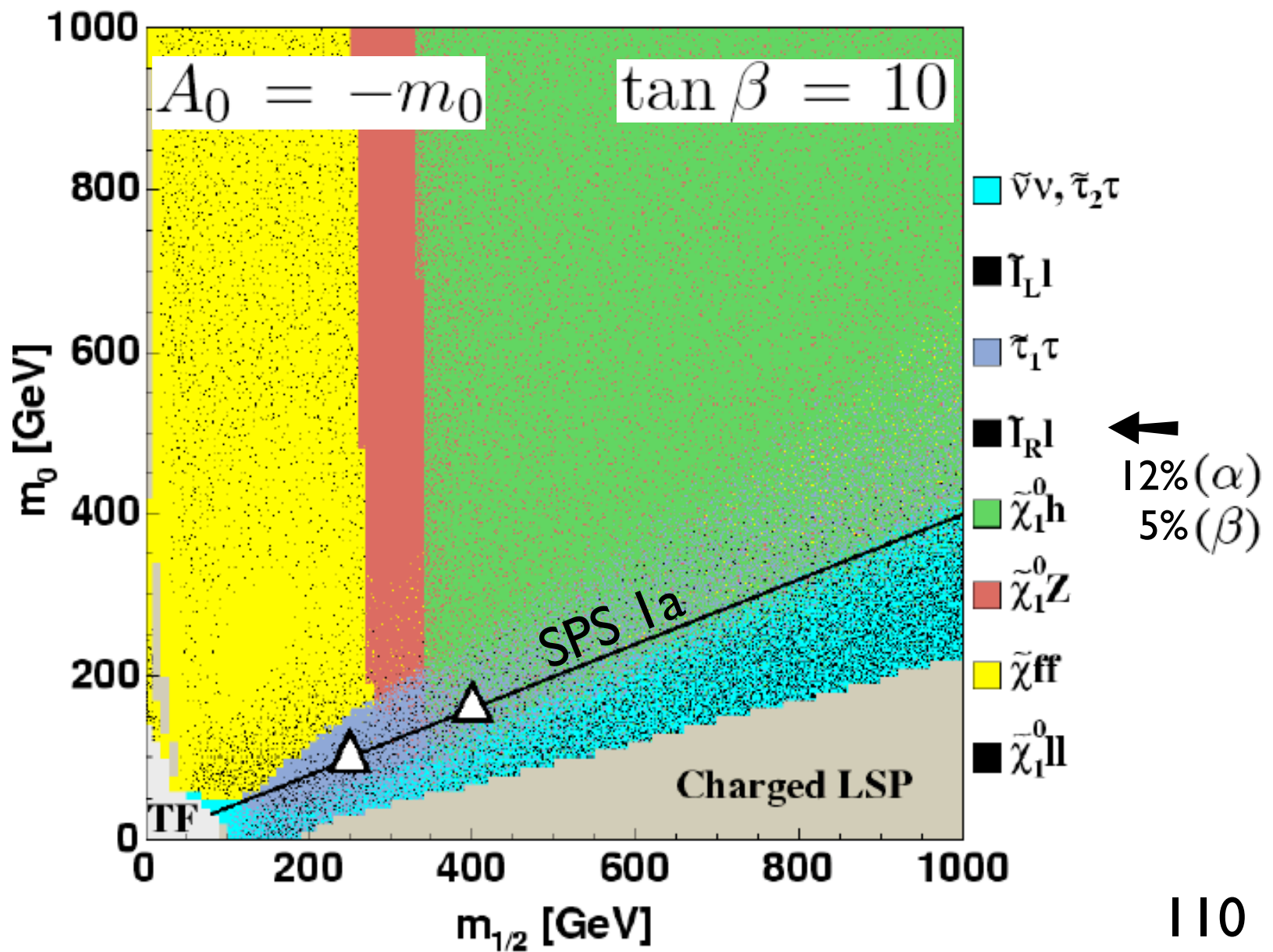
Squark Branching Ratios (\tilde{b}_1)



Squark Branching Ratios (\tilde{b}_2)



Neutralino Branching Ratios ($\tilde{\chi}_2^0$)



SPS Ia (line)

$$m_0 = -A_0 = 0.4 m_{1/2}$$
$$\tan \beta = 10, \quad \mu > 0$$

Two particular points on the line:

$$(\alpha) : \quad m_0 = 100 \text{ GeV}, \quad m_{1/2} = 250 \text{ GeV}$$

$$(\beta) : \quad m_0 = 160 \text{ GeV}, \quad m_{1/2} = 400 \text{ GeV}$$

Spectrum

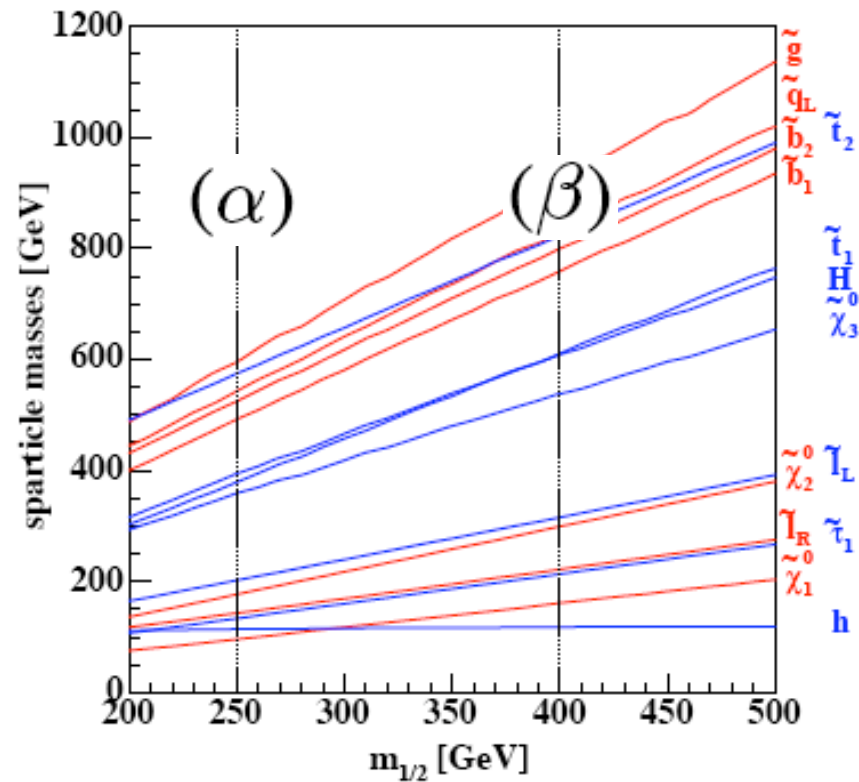
Point	\tilde{g}	\tilde{d}_L	\tilde{d}_R	\tilde{u}_L	\tilde{u}_R	\tilde{b}_2	\tilde{b}_1	\tilde{t}_2	\tilde{t}_1
(α)	595.2	543.0	520.1	537.2	520.5	524.6	491.9	574.6	379.1
(β)	915.5	830.1	799.5	826.3	797.3	800.2	759.4	823.8	610.4
	\tilde{e}_L	\tilde{e}_R	$\tilde{\tau}_2$	$\tilde{\tau}_1$	$\tilde{\nu}_{eL}$	$\tilde{\nu}_{\tau L}$		H^\pm	A
(α)	202.1	143.0	206.0	133.4	185.1	185.1		401.8	393.6
(β)	315.6	221.9	317.3	213.4	304.1	304.1		613.9	608.3
	$\tilde{\chi}_4^0$	$\tilde{\chi}_3^0$	$\tilde{\chi}_2^0$	$\tilde{\chi}_1^0$	$\tilde{\chi}_2^\pm$	$\tilde{\chi}_1^\pm$		H	h
(α)	377.8	358.8	176.8	96.1	378.2	176.4		394.2	114.0
(β)	553.3	538.4	299.1	161.0	553.3	299.0		608.9	117.9

as determined by ISASUSY 7.58 by integrating RGE's

in **bold**: particles used in study

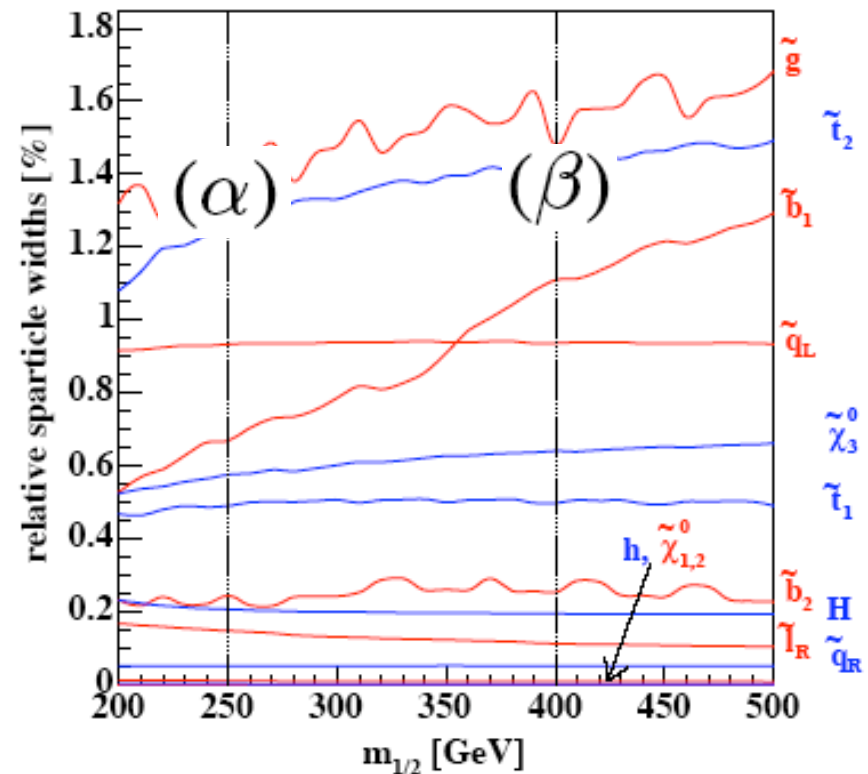
SPS Ia line

Masses



Widths

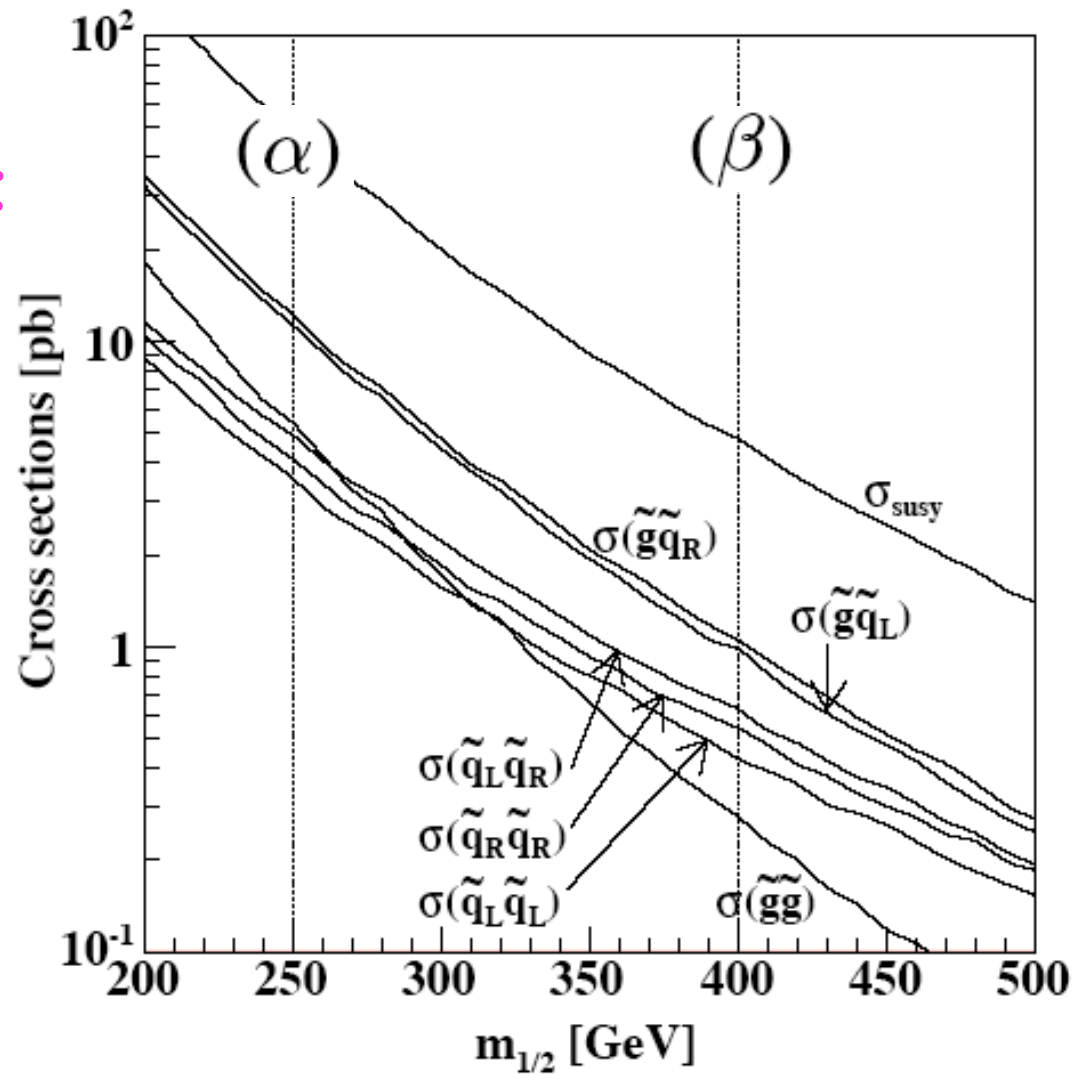
ISASUSY 7.58



$\sim 1\%$ of mass

Cross sections:

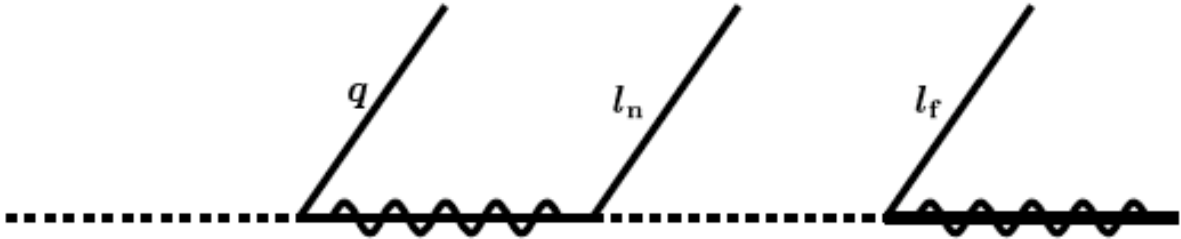
squark is often
produced together
with gluino,
or from gluino decay



[pb]

	$\sigma(\text{SUSY})$	$\sigma(\tilde{g}\tilde{g})$	$\sigma(\tilde{g}\tilde{q}_L)$	$\sigma(\tilde{g}\tilde{q}_R)$	$\sigma(\tilde{q}_L\tilde{q}_L)$	$\sigma(\tilde{q}_L\tilde{q}_R)$	$\sigma(\tilde{q}_R\tilde{q}_R)$
(α)	49.3	5.3	11.4	12.3	3.5	4.8	4.1
(β)	4.76	0.29	0.97	1.06	0.44	0.61	0.53

Quantifying the cascade:



(α)	$\sigma = 32.8 \text{ pb}$	\tilde{q}_L	31.4%	$\tilde{\chi}_2^0$ 12.1%	\tilde{l}_R 100%	$\tilde{\chi}_1^0$	1245 fb
	$\sigma = 7.7 \text{ pb}$	\tilde{b}_1	35.5%				329 fb
	$\sigma = 4.3 \text{ pb}$	\tilde{b}_2	18.0%				94 fb
							<hr/> 1669 fb
(β)	$\sigma = 3.21 \text{ pb}$	\tilde{q}_L	31.9%	$\tilde{\chi}_2^0$ 5.3%	\tilde{l}_R 100%	$\tilde{\chi}_1^0$	54.3 fb
	$\sigma = 0.50 \text{ pb}$	\tilde{b}_1	23.6%				6.3 fb
	$\sigma = 0.31 \text{ pb}$	\tilde{b}_2	8.8%				1.5 fb
							<hr/> 62.0 fb

Reduction (rate and BR) from (α) to (β)

Kinematics - lepton pair (simple example):

$$\tilde{\chi}_2^0 \rightarrow \tilde{l}_R l_n \quad \text{“n” = “near”}$$

$$m_{\tilde{\chi}_2^0}^2 = (p_{\tilde{l}_R} + p_{l_n})^2 = m_{\tilde{l}_R}^2 + 2p_{l_n} \cdot p_{\tilde{l}_R}$$

Rest frame of \tilde{l}_R :

$$|\mathbf{p}_{l_n}| = \frac{m_{\tilde{\chi}_2^0}^2 - m_{\tilde{l}_R}^2}{2m_{\tilde{l}_R}}$$

$$\tilde{l}_R \rightarrow \tilde{\chi}_1^0 l_f \quad \text{“f” = “far”}$$

$$m_{\tilde{\chi}_1^0}^2 = (p_{\tilde{l}_R} - p_{l_f})^2 = m_{\tilde{l}_R}^2 - 2p_{l_f} \cdot p_{\tilde{l}_R}$$

Rest frame of \tilde{l}_R :

$$|\mathbf{p}_{l_f}| = \frac{m_{\tilde{l}_R}^2 - m_{\tilde{\chi}_1^0}^2}{2m_{\tilde{l}_R}}$$

Maximum di-lepton mass:

Back-to-back in \tilde{l}_R Rest Frame:

$$(m_{ll}^{\max})^2 = 4|\mathbf{p}_{l_n}||\mathbf{p}_{l_f}| = \frac{(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{l}_R}^2)(m_{\tilde{l}_R}^2 - m_{\tilde{\chi}_1^0}^2)}{m_{\tilde{l}_R}^2}$$

Prototype of “endpoint formulas”

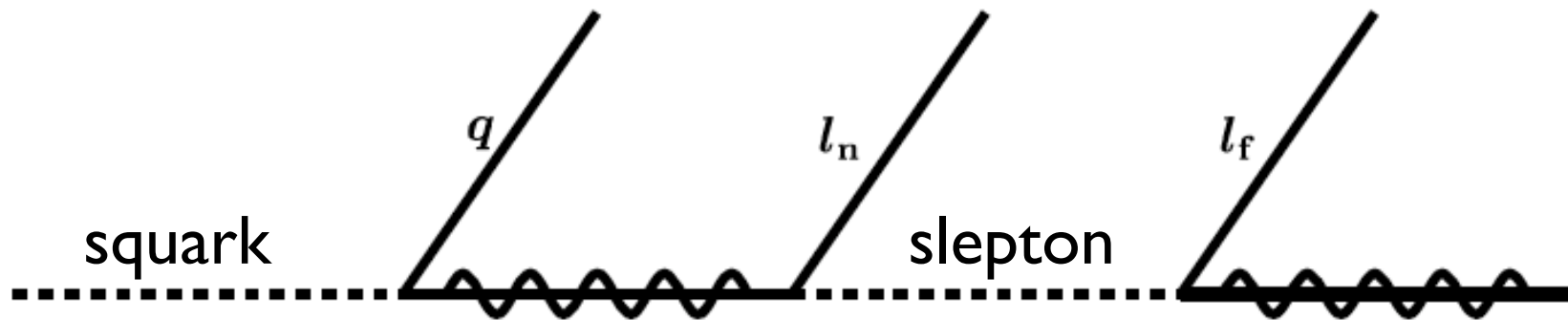
One kinematical endpt is related to various (3) masses of unstable particles:

$$m_{\tilde{\chi}_2^0} \quad m_{\tilde{l}_R} \quad m_{\tilde{\chi}_1^0}$$

Need more such formulas!

Add the squark:

$$\tilde{q}_L \rightarrow \tilde{\chi}_2^0 q \rightarrow \tilde{l}_R q l_n \rightarrow \tilde{\chi}_1^0 q l_n l_f$$



More invariants and endpoints:

$$m_{qll} \quad m_{ql_n} \quad m_{ql_f} \quad m_{ll}$$

Four endpoints and four masses:

$$m_{\tilde{q}_L} \quad m_{\tilde{\chi}_2^0} \quad m_{\tilde{l}_R} \quad m_{\tilde{\chi}_1^0}$$

Complication 1:

The two leptons can not be distinguished

For each event, form

$$m_{ql(\text{low})} < m_{ql(\text{high})} \quad \text{well defined}$$

Complication 2:

For some invariants, there are multiple cases:
endpoint formula depends on mass ratios

Complication 3:

Multiple squark masses; widths

Complication 4:

Endpoints not always linearly independent

B.C.Allanach et al, hep-ph/0007009 (conditions rephrased):

$$(m_{ll}^{\max})^2 = \frac{(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{l}_R}^2)(m_{\tilde{l}_R}^2 - m_{\tilde{\chi}_1^0}^2)}{m_{\tilde{l}_R}^2}$$

one case

mass ratios of adjacent
sparticles in chain

$$(m_{qll}^{\max})^2 = \left\{ \begin{array}{ll} \frac{(m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\chi}_1^0}^2)}{m_{\tilde{\chi}_2^0}^2} & \text{for } \frac{m_{\tilde{q}_L}}{m_{\tilde{\chi}_2^0}} > \frac{m_{\tilde{\chi}_2^0}}{m_{\tilde{l}_R}} \frac{m_{\tilde{l}_R}}{m_{\tilde{\chi}_1^0}} \quad (1) \\ \frac{(m_{\tilde{q}_L}^2 m_{\tilde{l}_R}^2 - m_{\tilde{\chi}_2^0}^2 m_{\tilde{\chi}_1^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{l}_R}^2)}{m_{\tilde{\chi}_2^0}^2 m_{\tilde{l}_R}^2} & \text{for } \frac{m_{\tilde{\chi}_2^0}}{m_{\tilde{l}_R}} > \frac{m_{\tilde{l}_R}}{m_{\tilde{\chi}_1^0}} \frac{m_{\tilde{q}_L}}{m_{\tilde{\chi}_2^0}} \quad (2) \\ \frac{(m_{\tilde{q}_L}^2 - m_{\tilde{l}_R}^2)(m_{\tilde{l}_R}^2 - m_{\tilde{\chi}_1^0}^2)}{m_{\tilde{l}_R}^2} & \text{for } \frac{m_{\tilde{l}_R}}{m_{\tilde{\chi}_1^0}} > \frac{m_{\tilde{q}_L}}{m_{\tilde{\chi}_2^0}} \frac{m_{\tilde{\chi}_2^0}}{m_{\tilde{l}_R}} \quad (3) \\ (m_{\tilde{q}_L} - m_{\tilde{\chi}_1^0})^2 & \text{otherwise} \quad (4) \end{array} \right\}$$

four cases

$$(m_{ql(\text{low})}^{\max}, m_{ql(\text{high})}^{\max}) = \left\{ \begin{array}{ll} (m_{ql_n}^{\max}, m_{ql_f}^{\max}) & \text{for } 2m_{\tilde{l}_R}^2 > m_{\tilde{\chi}_1^0}^2 + m_{\tilde{\chi}_2^0}^2 > 2m_{\tilde{\chi}_1^0} m_{\tilde{\chi}_2^0} \quad (1) \\ (m_{ql(\text{eq})}^{\max}, m_{ql_f}^{\max}) & \text{for } m_{\tilde{\chi}_1^0}^2 + m_{\tilde{\chi}_2^0}^2 > 2m_{\tilde{l}_R}^2 > 2m_{\tilde{\chi}_1^0} m_{\tilde{\chi}_2^0} \quad (2) \\ (m_{ql(\text{eq})}^{\max}, m_{ql_n}^{\max}) & \text{for } m_{\tilde{\chi}_1^0}^2 + m_{\tilde{\chi}_2^0}^2 > 2m_{\tilde{\chi}_1^0} m_{\tilde{\chi}_2^0} > 2m_{\tilde{l}_R}^2 \quad (3) \end{array} \right\}$$

three cases

where

$$(m_{ql_n}^{\max})^2 = \frac{(m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{l}_R}^2)}{m_{\tilde{\chi}_2^0}^2}$$

$$(m_{ql_i}^{\max})^2 = \frac{(m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{\tilde{l}_R}^2 - m_{\tilde{\chi}_1^0}^2)}{m_{\tilde{l}_R}^2}$$

$$(m_{ql(\text{eq})}^{\max})^2 = \frac{(m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{\tilde{l}_R}^2 - m_{\tilde{\chi}_1^0}^2)}{(2m_{\tilde{l}_R}^2 - m_{\tilde{\chi}_1^0}^2)}$$

Finally:

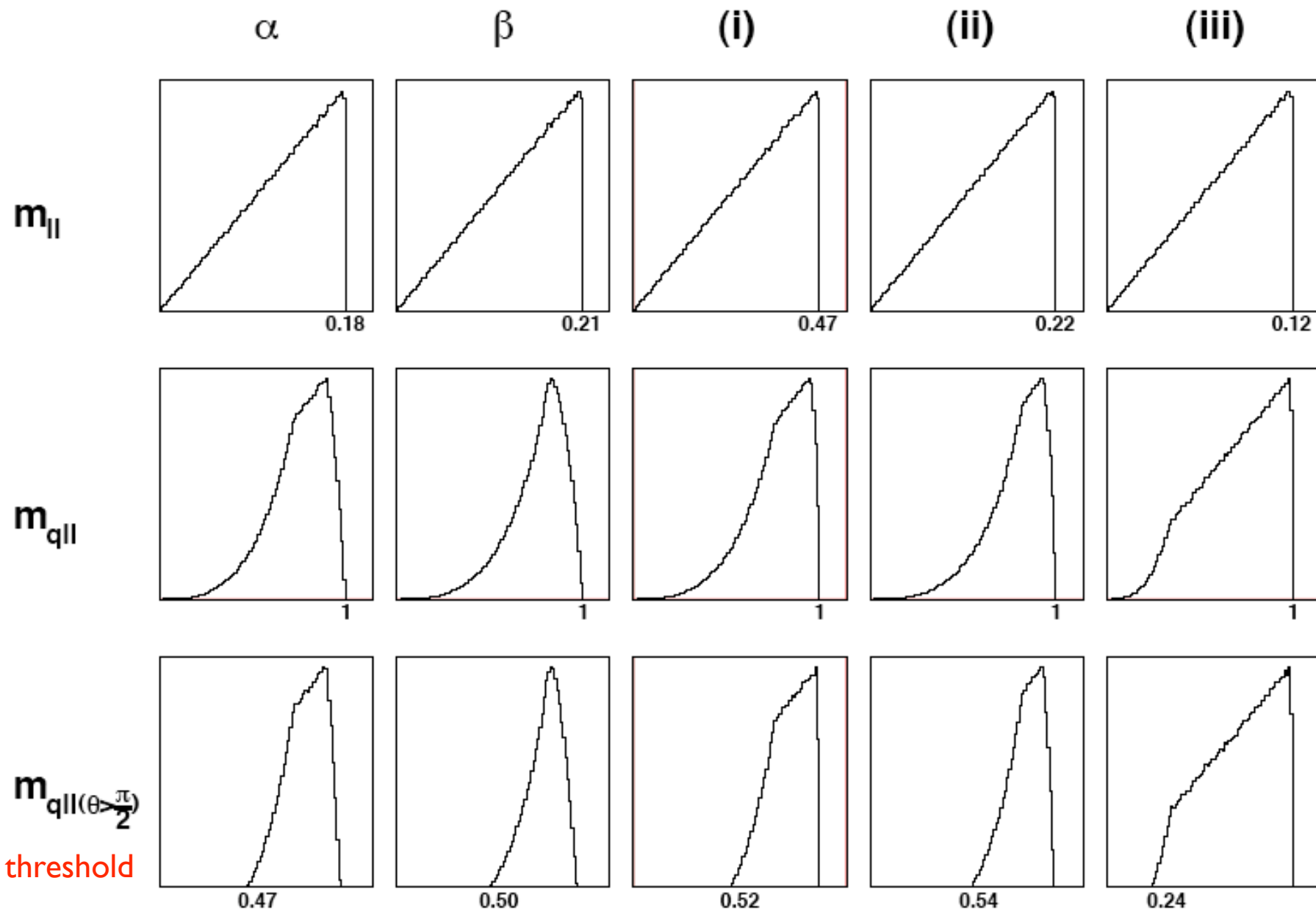
$$(m_{qll(\theta > \frac{\pi}{2})}^{\min})^2 = \left[(m_{\tilde{q}_L}^2 + m_{\tilde{\chi}_2^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{l}_R}^2)(m_{\tilde{l}_R}^2 - m_{\tilde{\chi}_1^0}^2) \right. \\ \left. - (m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2) \sqrt{(m_{\tilde{\chi}_2^0}^2 + m_{\tilde{l}_R}^2)^2 (m_{\tilde{l}_R}^2 + m_{\tilde{\chi}_1^0}^2)^2 - 16m_{\tilde{\chi}_2^0}^2 m_{\tilde{l}_R}^4 m_{\tilde{\chi}_1^0}^2} \right. \\ \left. + 2m_{\tilde{l}_R}^2 (m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\chi}_1^0}^2) \right] (4m_{\tilde{l}_R}^2 m_{\tilde{\chi}_2^0}^2)^{-1} \quad \text{one case}$$

θ is opening angle between leptons in \tilde{l}_R rest frame

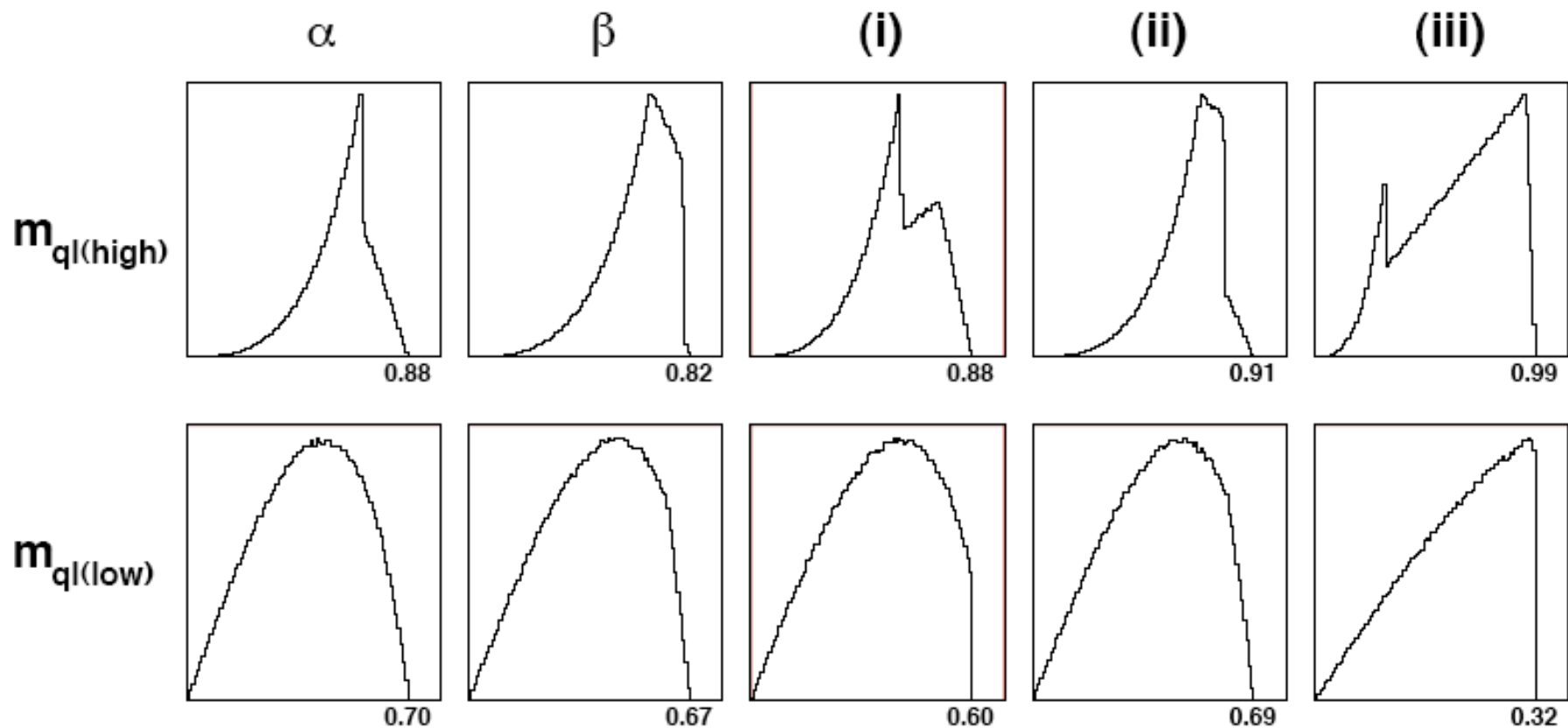
Over-all: 4×3 cases, denoted (1,1), (1,2), etc.

(9 of 12 are realized)

Theory curves (phase space)



Theory curves (phase space) cont.



α and β are on the SPS Ia line, other points “random”

Conclusions:
Many different shapes!
Edges not linear!

values in units of m_{ql}^{\max}

Inverting endpoint formulas

Endpoint formulas can be inverted

Complications:

nonlinear (rational, sqrt)

several cases (mass regions)

m_{ql}^{\max}

four cases

$m_{ql(\text{low})}^{\max}, m_{ql(\text{high})}^{\max}$ three cases

Example:

Region (1,1):

m_{ql}^{\max}

$m_{ql(\text{low})}^{\max}, m_{ql(\text{high})}^{\max}$

$$m_{\tilde{\chi}_1^0}^2 = \frac{(b^2 - d^2)(b^2 - c^2)}{(c^2 + d^2 - b^2)^2} a^2$$

$$m_{\tilde{l}_R}^2 = \frac{c^2(b^2 - c^2)}{(c^2 + d^2 - b^2)^2} a^2$$

$$m_{\tilde{\chi}_2^0}^2 = \frac{c^2 d^2}{(c^2 + d^2 - b^2)^2} a^2$$

$$m_{\tilde{q}_L}^2 = \frac{c^2 d^2}{(c^2 + d^2 - b^2)^2} (c^2 + d^2 - b^2 + a^2)$$

$$a = m_{ll}^{\max}, \quad b = m_{ql}^{\max}, \quad c = m_{ql(\text{low})}^{\max}, \quad d = m_{ql(\text{high})}^{\max}.$$

Inverting endpt formulas, cont

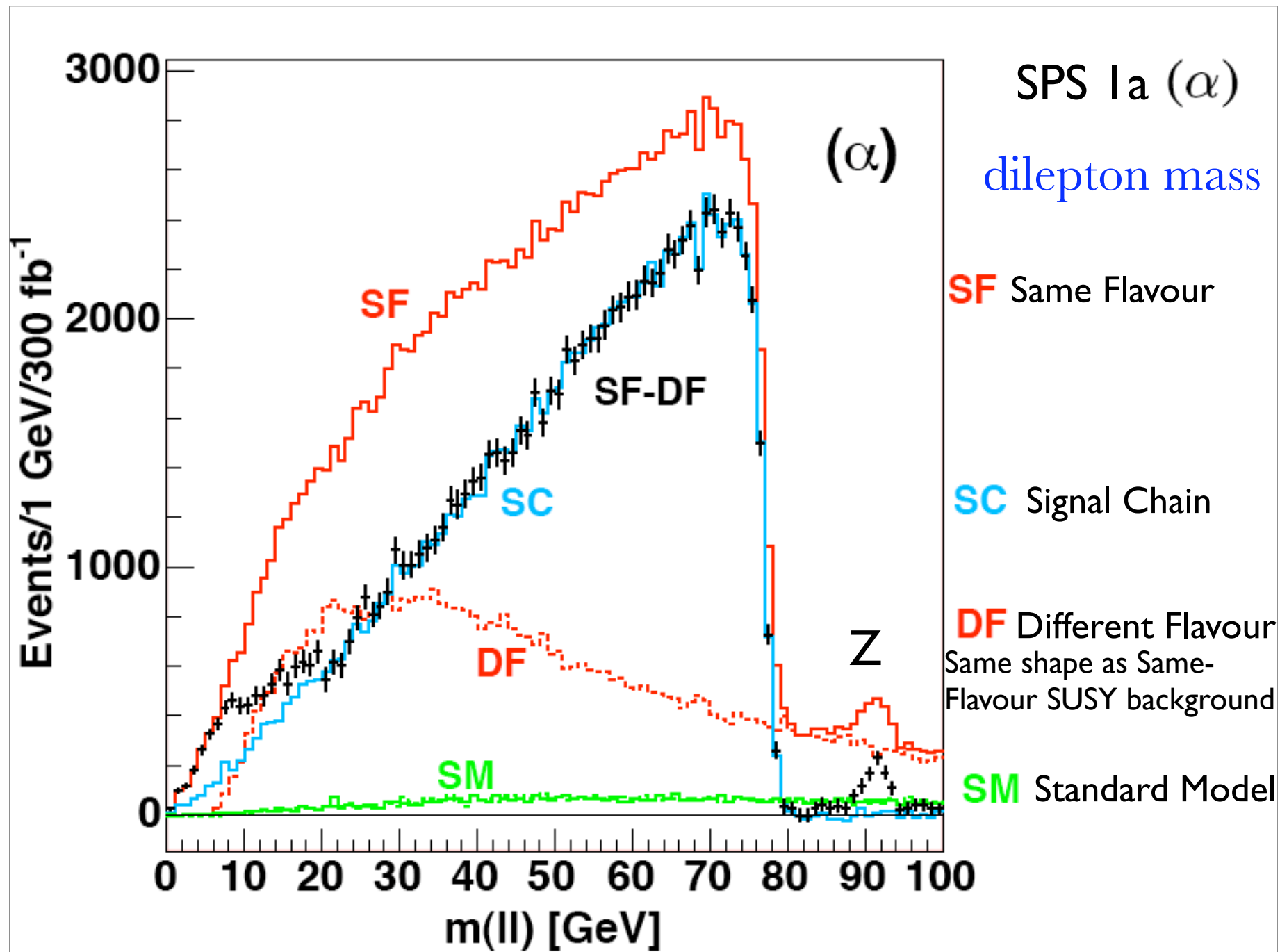
- If 4 endpts & 4 masses, (if linear) unique solution
- May have more endpts, system overconstrained
- Endpts have (different) uncertainties
- Use inversion formulas for start point of fit
- Composite formulas: multiple solutions!

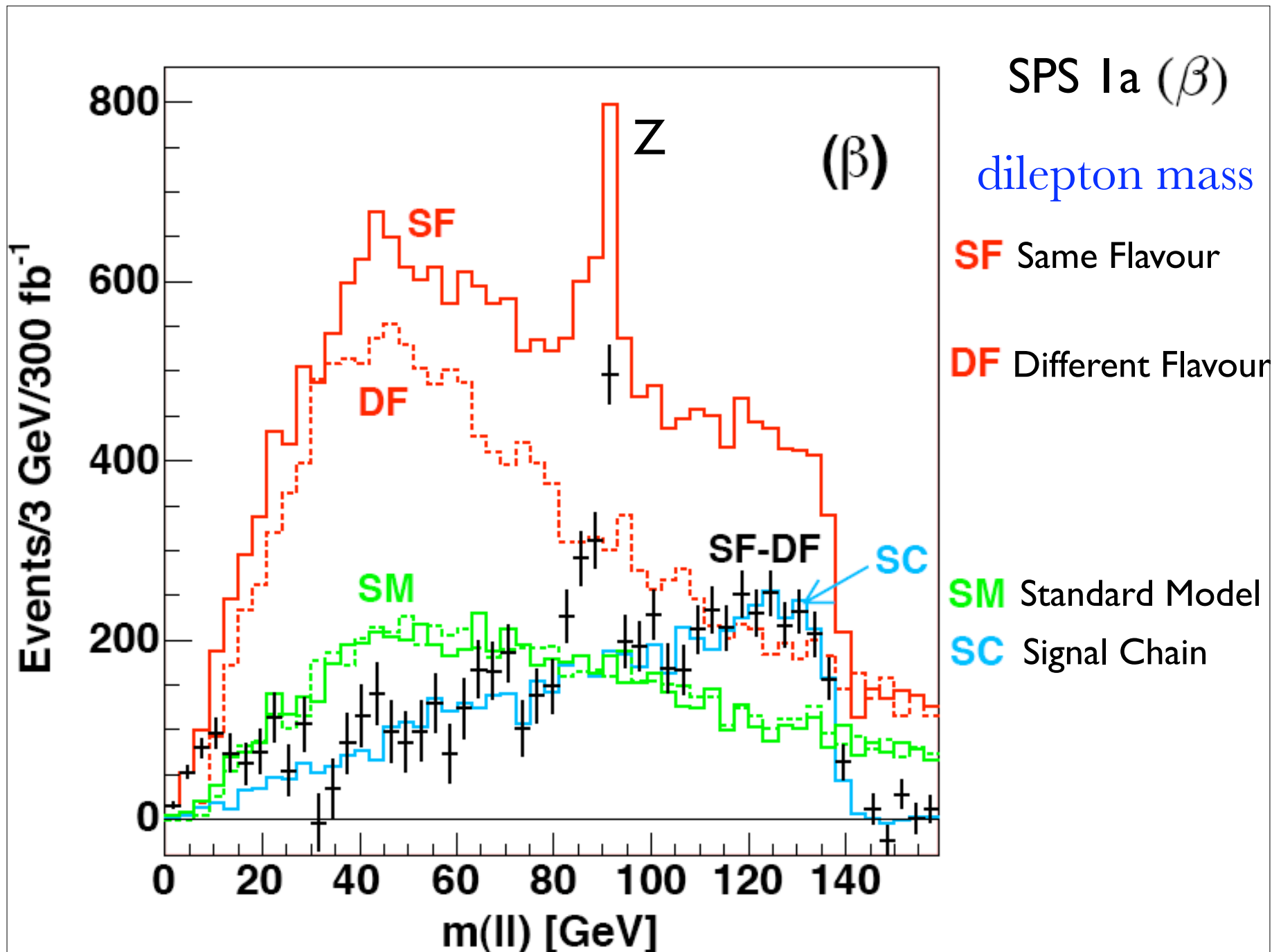
LHC simulation

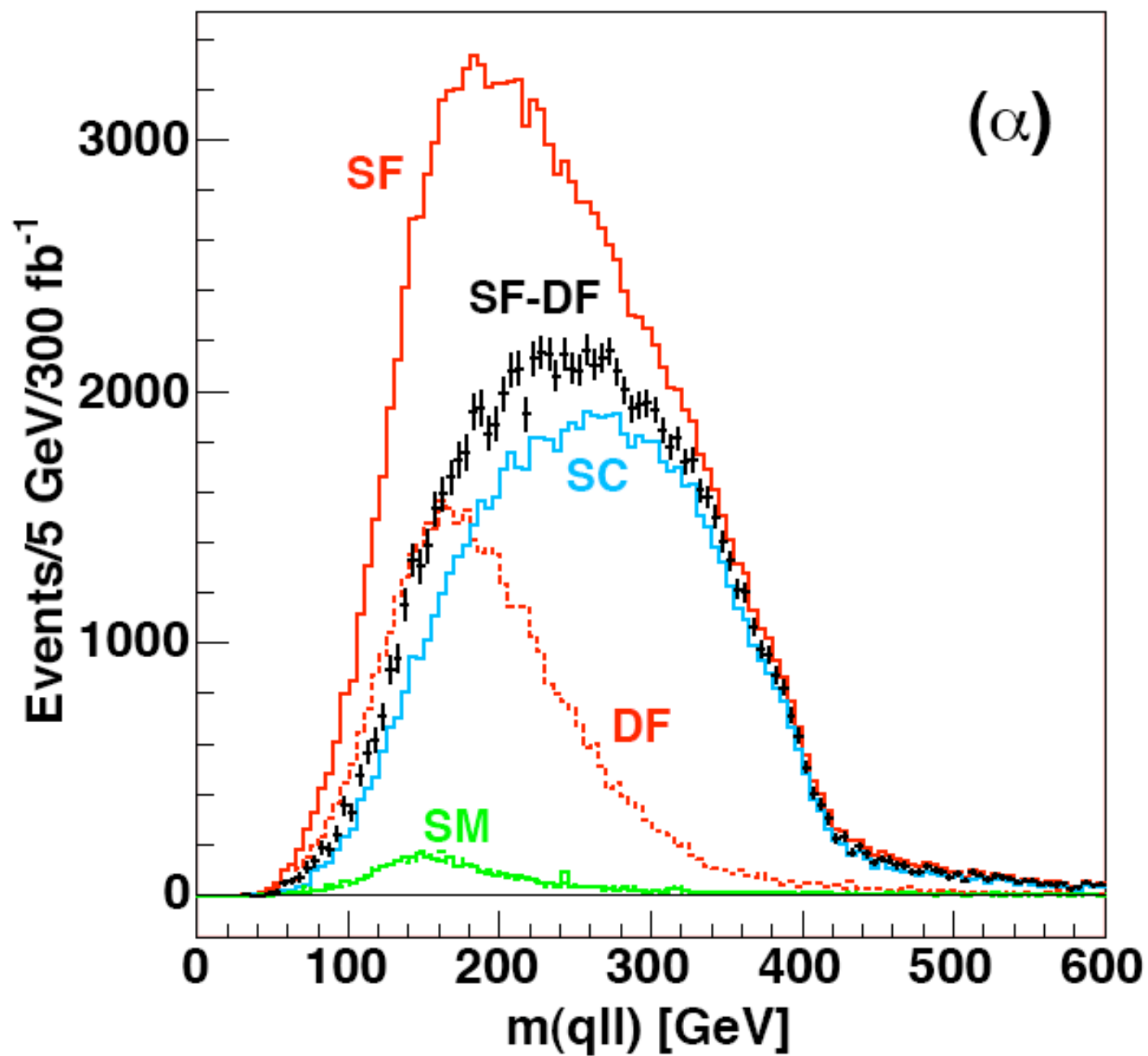
- ISAJET 7.58 defines low-energy model
- PYTHIA 6.2 with CTEQ 5L: Monte Carlo sample
- ATLFAST 2.60 simulates ATLAS detector
- precuts:
 - At least three jets, satisfying: $p_T^{\text{jet}} > 150, 100, 50 \text{ GeV}$
 - $E_{T,\text{miss}} > \max(100 \text{ GeV}, 0.2M_{\text{eff}})$
with $M_{\text{eff}} \equiv E_{T,\text{miss}} + \sum_{i=1}^3 p_{T,i}^{\text{jet}}$
 - Two isolated opposite-sign same-flavour leptons (e or μ),
satisfying $p_T^{\text{lep}} > 20, 10 \text{ GeV}$

SM background: 95% $t\bar{t}$

Aim: determine/study expected accuracy

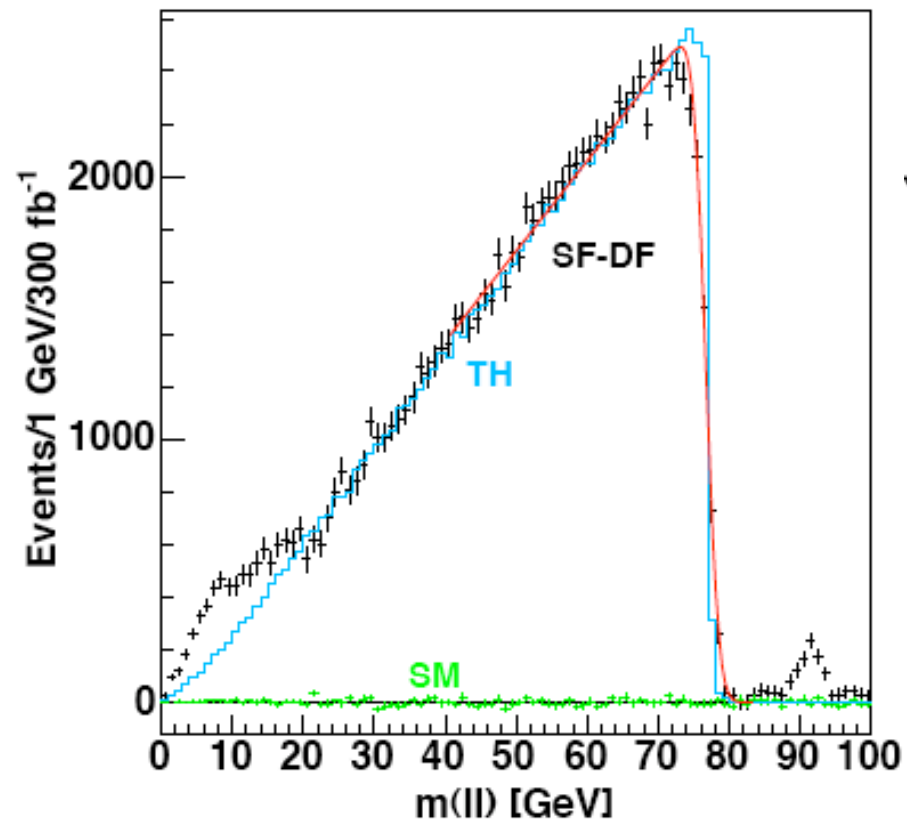




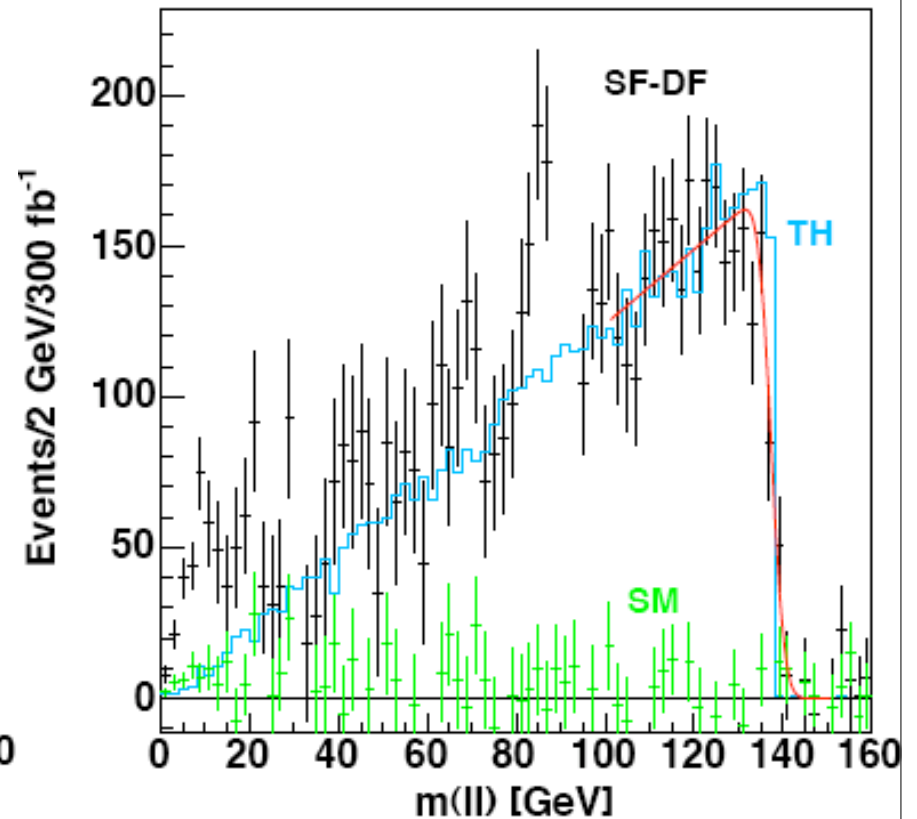


Fitting dilepton edge

(α)

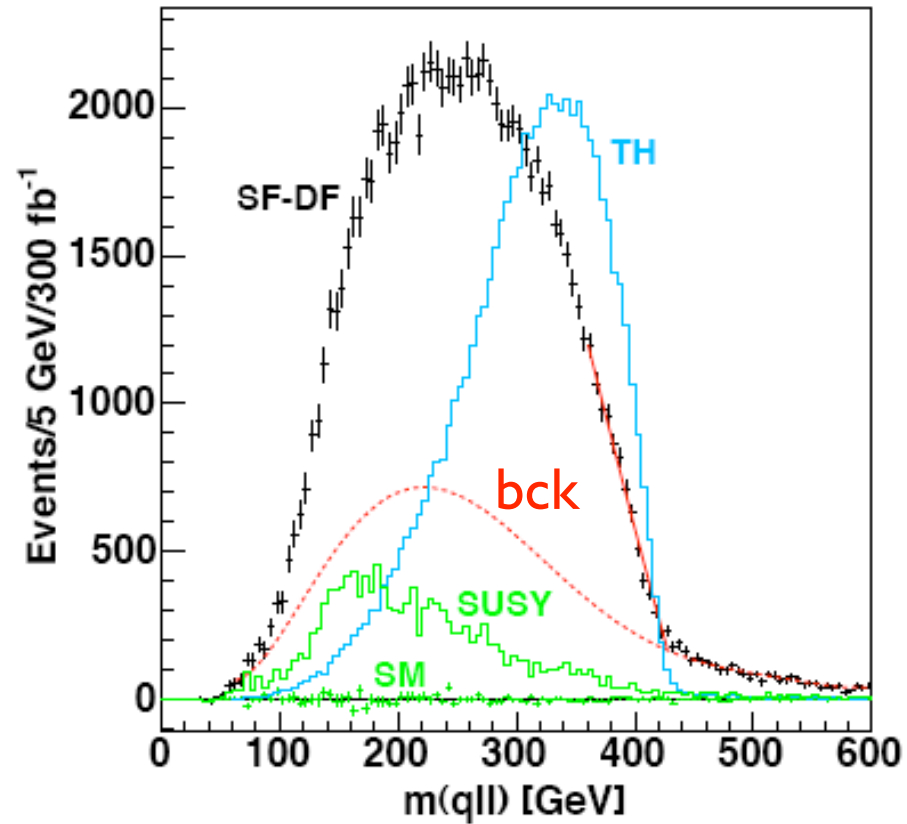


(β)

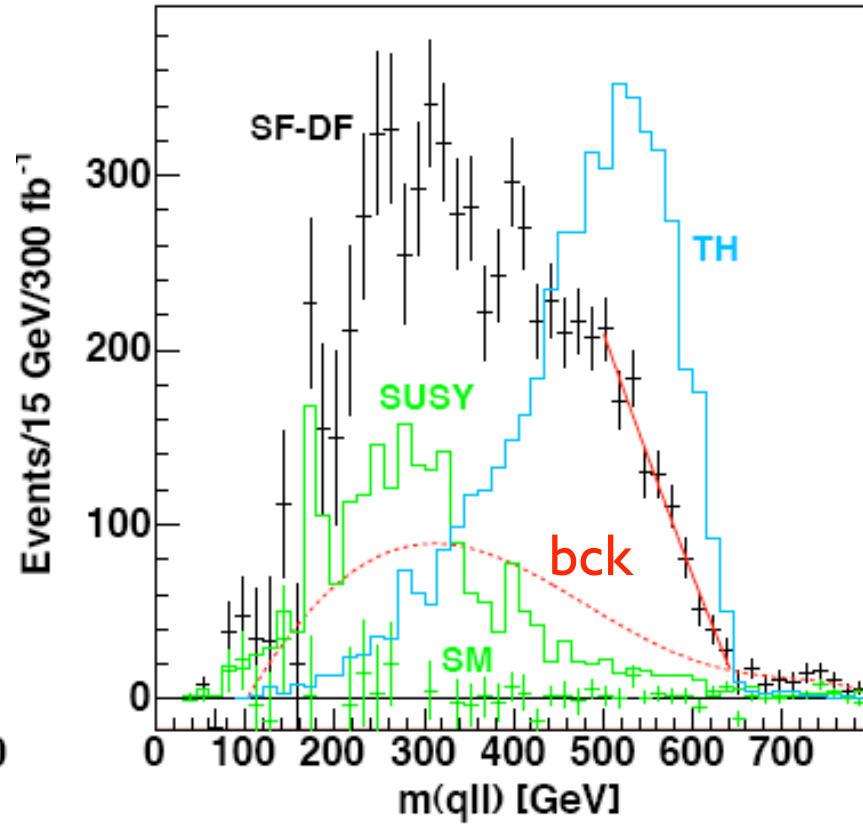


Fitting qll edge

(α)



(β)




background estimated from mixed events

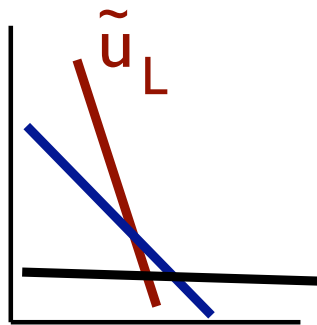
Fit results (endpoints) 300 fb⁻¹

Edge	Nominal Value [GeV]	Fit Value [GeV]	Energy Scale Error (σ^{scale}) [GeV]	Statistical Error (σ^{stat}) [GeV]	Syst. Fit Error [GeV]
(α)					
m_{ll}^{max}	77.07	76.72	0.08	0.04	0.1
m_{qll}^{max}	425.9	427.7	2.1	0.9	0.5
$m_{ql(\text{low})}^{\text{max}}$	298.5	300.7	1.5	0.9	0.5
$m_{ql(\text{high})}^{\text{max}}$	375.8	374.0	1.9	1.0	0.5
$m_{qll(\theta > \frac{\pi}{2})}^{\text{min}}$	200.7	-	1.0	2.2	2.0
$m_{bll(\theta > \frac{\pi}{2})}^{\text{min}}$	183.1	-	0.9	4.5	4.0
(β)					
m_{ll}^{max}	137.9	137.4	0.14	0.5	0.1
m_{qll}^{max}	649.1	647.0	3.2	5.0	3.0
$m_{ql(\text{low})}^{\text{max}}$	436.6	443.0	2.2	6.3	4.0
$m_{ql(\text{high})}^{\text{max}}$	529.9	520.5	2.6	5.5	3.0
$m_{qll(\theta > \frac{\pi}{2})}^{\text{min}}$	325.7	-	1.6	13.0	10.0

Multiple squark masses:

	\tilde{b}_1	\tilde{b}_2	\tilde{u}_L	\tilde{d}_L
SPS Ia (α)	491.9	524.6	537.2	543.0
SPS Ia (β)	759.4	800.2	826.3	830.1


two endpoints



Higher mass: higher endpoint
but lower rate - background -
blurring of endpoint

Extraction of masses

- simulate 10,000 ATLAS ‘experiments’
- focus on statistical uncertainty
- each endpoint: gaussian distribution
- invert endpoint formulas, fit masses
- what is chance of finding correct minimum?

Following Allanach et al, each endpt E_i^{exp} taken as:

$$E_i^{\text{exp}} = E_i^{\text{nom}} + A_i \sigma_i^{\text{stat}} + B \sigma_i^{\text{scale}}$$

A, B picked from gaussian distribution, mean 0, width 1

One A for each endpoint,

one B for m_{ll} , other B for endpoints involving jets

Minimize:

$$\Sigma = [\mathbf{E}^{\text{exp}} - \mathbf{E}^{\text{th}}(\mathbf{m})]^T \mathbf{W} [\mathbf{E}^{\text{exp}} - \mathbf{E}^{\text{th}}(\mathbf{m})]$$

\mathbf{W} inverse error/correlation matrix

determine masses



SPS Ia (α) $\Delta\Sigma \leq 1$

nominal

correct fit

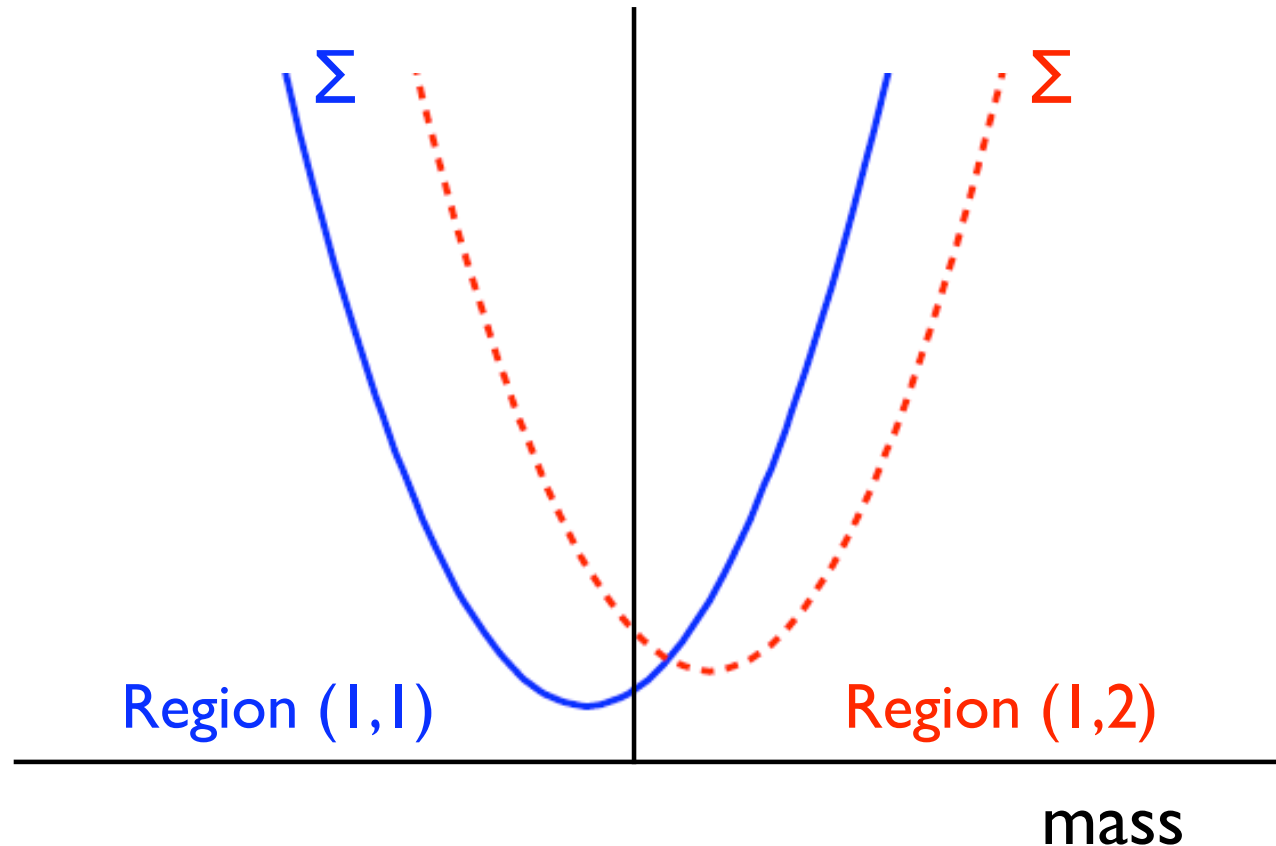
false fit

	Nom	$(1,1)$			$(1,2)$		
		$\langle m \rangle$	σ	γ_1	$\langle m \rangle$	σ	γ_1
$m_{\tilde{\chi}_1^0}$	96.1	96.3	3.8	0.2	85.3	3.4	0.1
$m_{\tilde{l}_R}$	143.0	143.2	3.8	0.2	130.4	3.7	0.1
$m_{\tilde{\chi}_2^0}$	176.8	177.0	3.7	0.2	165.5	3.4	0.1
$m_{\tilde{q}_L}$	537.2	537.5	6.1	0.1	523.2	5.1	0.1
$m_{\tilde{b}_1}$	491.9	492.4	13.4	0.0	469.6	13.3	0.1
$m_{\tilde{l}_R} - m_{\tilde{\chi}_1^0}$	46.9	46.9	0.3	0.0	45.1	0.7	-0.2
$m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$	80.8	80.8	0.2	0.0	80.2	0.3	-0.1
$m_{\tilde{q}_L} - m_{\tilde{\chi}_1^0}$	441.2	441.3	3.1	0.0	438.0	2.7	0.0
$m_{\tilde{b}_1} - m_{\tilde{\chi}_1^0}$	395.9	396.2	12.0	0.0	384.4	12.0	0.1

Note: Three lightest masses are very correlated

Problem due to compositeness of formulas:

If masses are close to border of 'region', may find a similar-quality or better minimum in 'other' region



How likely is a **false** minimum?

Depends on cut $\Delta\Sigma$ (level of confidence)

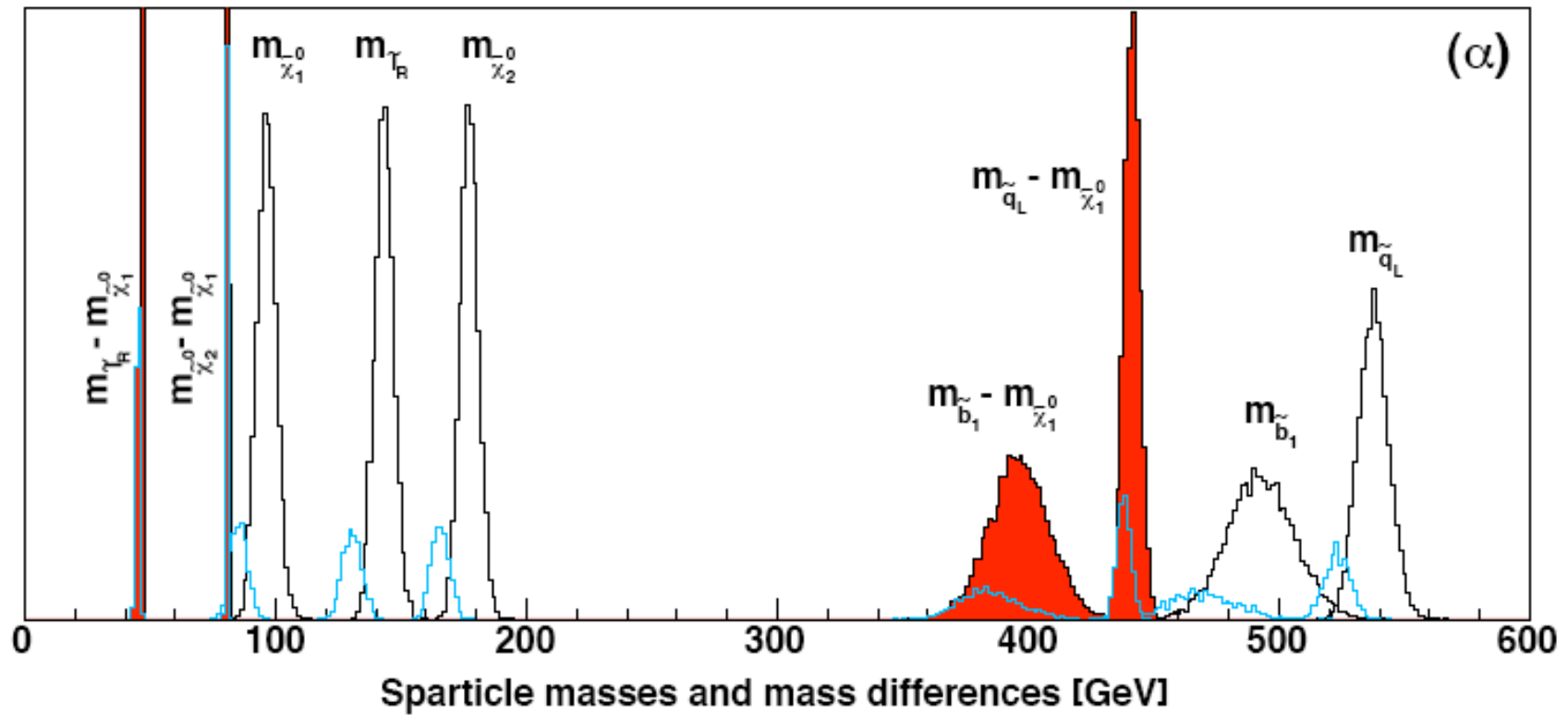
SPS Ia (α)

	# Minima	(1,1)	(1,2)
$\Delta\Sigma \leq 0$	1.00	90%	10%
$\Delta\Sigma \leq 1$	1.12	94%	17%
$\Delta\Sigma \leq 3$	1.30	97%	33%
$\Delta\Sigma \leq \infty$	1.88	99%	88%

Example: $\Delta\Sigma \leq 3$

30% chance of finding *two* minima

Spread

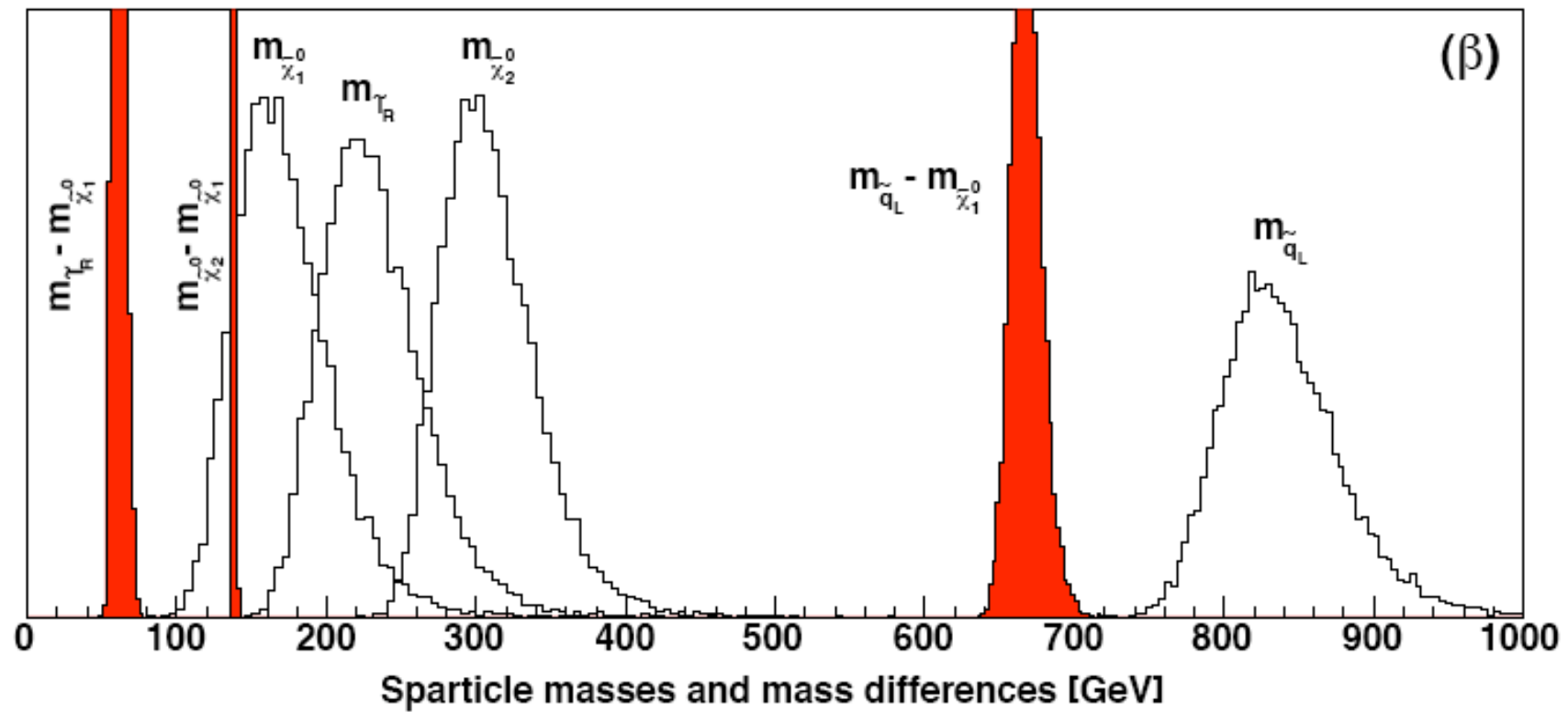


black: correct solution

red: mass differences

blue-green: false solution (area prop to probability)

Spread



false solutions not shown

LC input ("fixing" LSP mass)

SPS Ia (α)

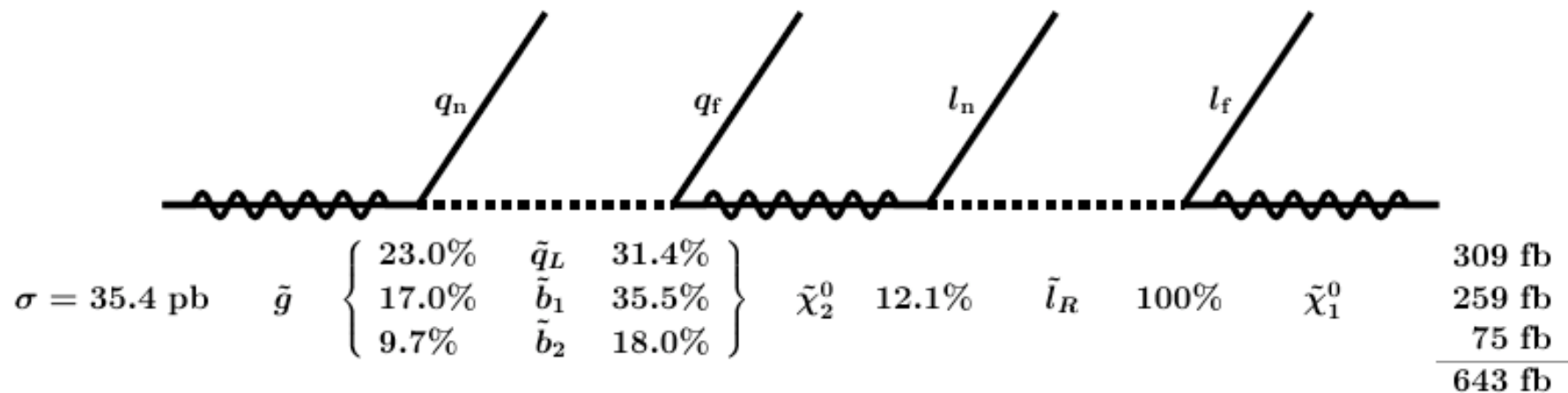
	Nom	$(1,1)$ $\langle m \rangle$	σ
$\tilde{\chi}_1^0$	96.05	96.05	0.05
\tilde{l}_R	142.97	142.97	0.29
$\tilde{\chi}_2^0$	176.82	176.82	0.17
\tilde{q}_L	537.25	537.2	2.5
\tilde{b}_1	491.92	492.1	11.7

Masses in GeV

SPS Ia (β)

		1 solution		2 solutions			
		$(1,2)/(1,3)/B$		$(1,2)$		$(1,3)$	
	Nom	$\langle m \rangle$	σ	$\langle m \rangle$	σ	$\langle m \rangle$	σ
$\tilde{\chi}_1^0$	161.02	161.02	0.05	161.02	0.05	161.02	0.05
\tilde{l}_R	221.86	221.15	3.26	222.22	1.32	217.48	1.01
$\tilde{\chi}_2^0$	299.05	299.15	0.57	299.11	0.53	299.05	0.52
\tilde{q}_L	826.29	826.1	6.3	825.9	5.8	828.6	5.5

Gluino cascade chain



SPS Ia numbers

Several new kinematical edges involving q_n

Only *one* new mass, need (minimum) only *one* more edge

Summary (SPS 1a)

- SPS 1a SUSY masses can be determined with precision 4-10 GeV
- Non-zero probability of fitting wrong minimum
- LC input on LSP mass removes ambiguity and increases precision significantly
- Gluino mass can be obtained using two b jets
- Better understanding of backgrounds very beneficial
- Understand better systematics of fitting (detector & theory)