

Inflation: the present status

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Present status of inflation

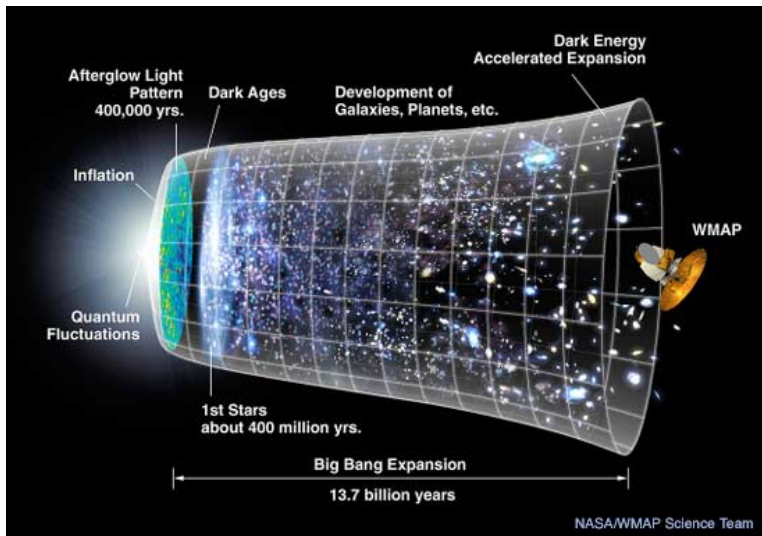
The simplest one-parametric inflationary models

Inflation and its smooth reconstruction in GR

Inflation and its smooth reconstruction in $f(R)$ gravity

Small local features in the power spectrum

Conclusions



Four epochs of the history of the Universe

$H \equiv \frac{\dot{a}}{a}$ where $a(t)$ is a scale factor of an isotropic homogeneous spatially flat universe (a Friedmann-Lemaître-Robertson-Walker background):

$$ds^2 = dt^2 - a^2(t)(dx^2 + dy^2 + dz^2) + \textit{small perturbations}$$

The history of the Universe in one line: four main epochs

$$? \longrightarrow DS \Longrightarrow FLWRD \Longrightarrow FLRWMD \Longrightarrow \overline{DS} \longrightarrow ?$$

Geometry

$$|\dot{H}| \ll H^2 \Longrightarrow H = \frac{1}{2t} \Longrightarrow H = \frac{2}{3t} \Longrightarrow |\dot{H}| \ll H^2$$

Physics

$$p \approx -\rho \Longrightarrow p = \rho/3 \Longrightarrow p \ll \rho \Longrightarrow p \approx -\rho$$

Duration in terms of the number of e-folds $\ln(a_{fin}/a_{in})$

> 60

~ 55

7.5

0.5

Main epochs of the Universe evolution – before 1979

The history of the Universe in one line: two main epochs

? \longrightarrow *FLWRD* \Longrightarrow *FLRWMD* \longrightarrow ?

Geometry

$$H = \frac{1}{2t} \Longrightarrow H = \frac{2}{3t}$$

Physics

$$p = \rho/3 \Longrightarrow p \ll \rho$$

Inflation

The inflationary scenario is based on the two cornerstone independent ideas (hypothesis):

1. Existence of **inflation** (or, quasi-de Sitter stage) – a stage of accelerated, close to exponential expansion of our Universe in the past preceding the hot Big Bang with decelerated, power-law expansion.
2. The origin of all inhomogeneities in the present Universe is the effect of **gravitational creation of particles and field fluctuations** during inflation from the adiabatic vacuum (no-particle) state for Fourier modes covering all observable range of scales (and possibly somewhat beyond).

Existing analogies in other areas of physics.

1. The present dark energy.
2. Creation of electrons and positrons in an external electric field.

Outcome of inflation

In the super-Hubble regime ($k \ll aH$) in the coordinate representation:

$$ds^2 = dt^2 - a^2(t)(\delta_{lm} + h_{lm})dx^l dx^m, \quad l, m = 1, 2, 3$$

$$h_{lm} = 2\zeta(\mathbf{r})\delta_{lm} + \sum_{a=1}^2 g^{(a)}(\mathbf{r}) e_{lm}^{(a)}$$

$$e_l^{l(a)} = 0, \quad g_{,l}^{(a)} e_m^{l(a)} = 0, \quad e_{lm}^{(a)} e^{lm(a)} = 1$$

ζ describes primordial scalar perturbations, g – primordial tensor perturbations (primordial gravitational waves (GW)).

The most important quantities:

$$n_s(k) - 1 \equiv \frac{d \ln P_\zeta(k)}{d \ln k}, \quad r(k) \equiv \frac{P_g}{P_\zeta}$$

In fact, metric perturbations h_{lm} are quantum (operators in the Heisenberg representation) and remain quantum up to the present time. But, after omitting of a very small part, decaying with time, they become commuting and, thus, equivalent to classical (c-number) stochastic quantities with the Gaussian statistics (up to small terms quadratic in ζ , g).

In particular:

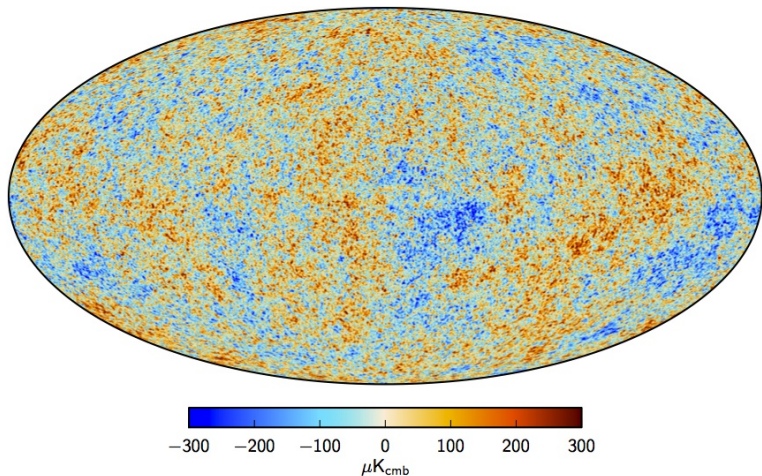
$$\hat{\zeta}_k = \zeta_k i(\hat{a}_k - \hat{a}_k^\dagger) + \mathcal{O}\left((\hat{a}_k - \hat{a}_k^\dagger)^2\right) + \dots + \mathcal{O}(10^{-100})(\hat{a}_k + \hat{a}_k^\dagger) + \dots,$$

The last term is time dependent, it is affected by physical decoherence and may become larger, but not as large as the second term.

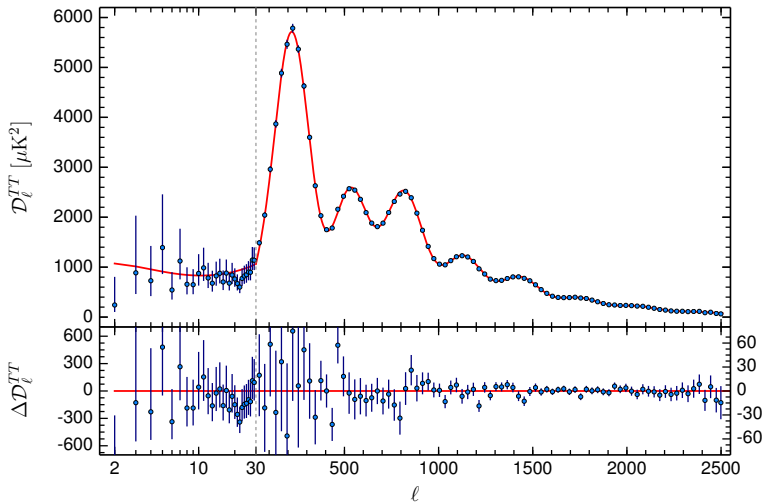
Remaining quantum coherence: deterministic correlation between \mathbf{k} and $-\mathbf{k}$ modes - shows itself in the appearance of acoustic oscillations (primordial oscillations in case of GW).

CMB temperature anisotropy

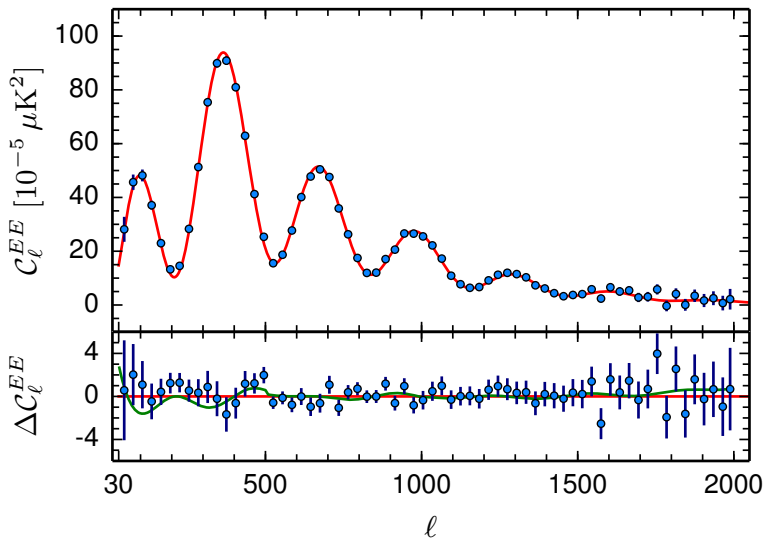
Planck-2015: P. A. R. Ade et al., arXiv:1502.01589



CMB temperature anisotropy multipoles



CMB E-mode polarization multipoles



Present status of inflation

Now we have quantitative observational data: the primordial spectrum of scalar perturbations has been measured and its deviation from the flat spectrum $n_s = 1$ in the first order in $|n_s - 1| \sim N_H^{-1}$ has been discovered (using the multipole range $\ell > 40$).

$$\langle \zeta^2(\mathbf{r}) \rangle = \int \frac{P_\zeta(k)}{k} dk, \quad P_\zeta(k) = (2.21^{+0.07}_{-0.08}) 10^{-9} \left(\frac{k}{k_0} \right)^{n_s-1}$$

$$k_0 = 0.05 \text{Mpc}^{-1}, \quad n_s - 1 = -0.035 \pm 0.005$$

Two fundamental observational constants of cosmology in addition to the three known ones (baryon-to-photon ratio, baryon-to-matter density and the cosmological constant). Existing inflationary models can predict (and predicted, in fact) one of them, namely $n_s - 1$, relating it finally to $N_H = \ln \frac{k_B T_\gamma}{\hbar H_0} \approx 67.2$ (note that $(1 - n_s)N_H \sim 2$).

From "proving" inflation to using it as a tool

Present status of inflation: transition from "proving" it in general and testing some of its simplest models to applying the inflationary paradigm to investigate particle physics at super-high energies and the actual history of the Universe in the remote past using real observational data on $n_s(k) - 1$ and $r(k)$.

The reconstruction approach – determining curvature and inflaton potential from observational data – a kind of inverse dynamical problem.

The most important quantities:

- 1) for classical gravity – H, \dot{H}
- 2) for super-high energy particle physics – m_{infl}^2 .

Physical scales related to inflation

"Naive" estimate where I use the reduced Planck mass

$$\tilde{M}_{Pl} = (8\pi G)^{-1/2}.$$

I. Curvature scale

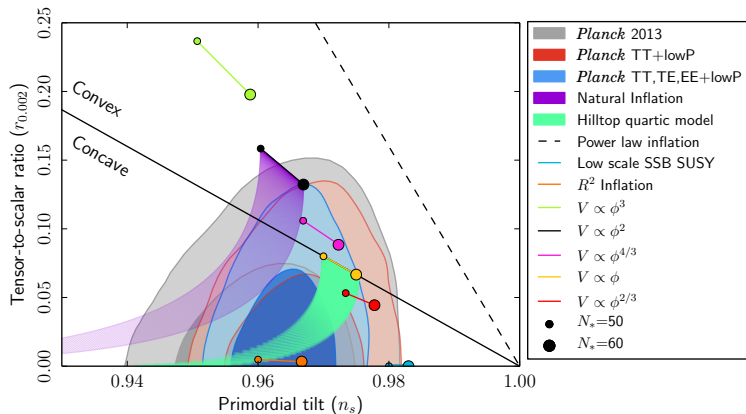
$$H \sim \sqrt{P_\zeta} \tilde{M}_{Pl} \sim 10^{14} \text{GeV}$$

II. Inflaton mass scale

$$|m_{infl}| \sim H \sqrt{|1 - n_s|} \sim 10^{13} \text{GeV}$$

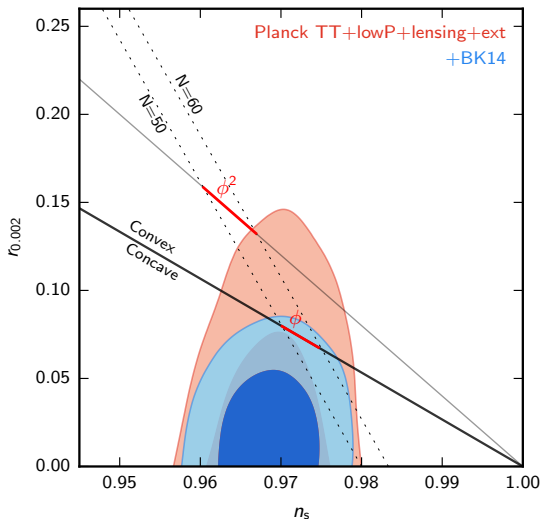
New range of mass scales significantly less than the GUT scale.

Direct approach: comparison with simple smooth models



Combined BICEP2/Keck Array/Planck results

P. A. R. Ade et al., Phys. Rev. Lett. 116, 031302 (2016)



The simplest models producing the observed scalar slope

$$\mathcal{L} = \frac{f(R)}{16\pi G}, \quad f(R) = R + \frac{R^2}{6M^2}$$

$$M = 2.6 \times 10^{-6} \left(\frac{55}{N} \right) M_{Pl} \approx 3.2 \times 10^{13} \text{ GeV}$$

$$n_s - 1 = -\frac{2}{N} \approx -0.036, \quad r = \frac{12}{N^2} \approx 0.004$$

$$N = \ln \frac{k_f}{k} = \ln \frac{T_\gamma}{k} - \mathcal{O}(10), \quad H_{dS}(N = 55) = 1.4 \times 10^{14} \text{ GeV}$$

The same prediction from a scalar field model with $V(\phi) = \frac{\lambda\phi^4}{4}$ at large ϕ and strong non-minimal coupling to gravity $\xi R\phi^2$ with $\xi < 0$, $|\xi| \gg 1$, including the Brout-Englert-Higgs inflationary model.

The simplest purely geometrical inflationary model

$$\begin{aligned}\mathcal{L} &= \frac{R}{16\pi G} + \frac{N^2}{288\pi^2 P_\zeta(k)} R^2 + (\text{small rad. corr.}) \\ &= \frac{R}{16\pi G} + 5 \times 10^8 R^2 + (\text{small rad. corr.})\end{aligned}$$

The quantum effect of creation of particles and field fluctuations works **twice** in this model:

- a) at super-Hubble scales during inflation, to generate space-time metric fluctuations;
- b) at small scales after inflation, to provide scalaron decay into pairs of matter particles and antiparticles (AS, 1980, 1981).

The most effective decay channel: into minimally coupled scalars with $m \ll M$. Then the formula

$$\frac{1}{\sqrt{-g}} \frac{d}{dt} (\sqrt{-g} n_s) = \frac{R^2}{576\pi}$$

(Ya. B. Zeldovich and A. A. Starobinsky, JETP Lett. 26, 252 (1977)) can be used for simplicity, but the full integral-differential system of equations for the Bogoliubov α_k, β_k coefficients and the average EMT was in fact solved in AS (1981). Sclaron decay into graviton pairs is suppressed (A. A. Starobinsky, JETP Lett. 34, 438 (1981)).

Possible microscopic origins of this phenomenological model.

1. Follow the purely geometrical approach and consider it as the specific case of the fourth order gravity in 4D

$$\mathcal{L} = \frac{R}{16\pi G} + AR^2 + BC_{\alpha\beta\gamma\delta}C^{\alpha\beta\gamma\delta} + (\text{small rad. corr.})$$

for which $A \gg 1$, $A \gg |B|$. Approximate scale (dilaton) invariance and absence of ghosts in the curvature regime $A^{-2} \ll (RR)/M_P^4 \ll B^{-2}$.

One-loop quantum-gravitational corrections are small (their imaginary parts are just the predicted spectra of scalar and tensor perturbations), non-local and qualitatively have the same structure modulo logarithmic dependence on curvature.

2. Another, completely different way:

consider the $R + R^2$ model as an **approximate** description of GR + a non-minimally coupled scalar field with a large negative coupling ξ ($\xi_{\text{conf}} = \frac{1}{6}$) in the gravity sector::

$$L = \frac{R}{16\pi G} - \frac{\xi R \phi^2}{2} + \frac{1}{2} \phi_{,\mu} \phi^{,\mu} - V(\phi), \quad \xi < 0, \quad |\xi| \gg 1 .$$

Geometrization of the scalar:

for a generic family of solutions during inflation and even for some period of non-linear scalar field oscillations after it, the scalar kinetic term can be neglected, so

$$\xi R \phi = -V'(\phi) + \mathcal{O}(|\xi|^{-1}) .$$

No conformal transformation, we remain in the the physical (Jordan) frame!

These solutions are the same as for $f(R)$ gravity with

$$L = \frac{f(R)}{16\pi G}, \quad f(R) = R - \frac{\xi R \phi^2(R)}{2} - V(\phi(R)).$$

For $V(\phi) = \frac{\lambda(\phi^2 - \phi_0^2)^2}{4}$, this just produces

$$f(R) = \frac{1}{16\pi G} \left(R + \frac{R^2}{6M^2} \right) \text{ with } M^2 = \lambda/24\pi\xi^2 G \text{ and } \phi^2 = |\xi|R/\lambda.$$

The same theorem is valid for a multi-component scalar field, as well as for the mixed Higgs- R^2 model (M. He, A. A. Starobinsky, J. Yokoyama, arXiv:1804.00409).

More generally, R^2 inflation (with an arbitrary n_s, r) serves as an intermediate **dynamical** attractor for a large class of scalar-tensor gravity models.

Inflation in GR

Inflation in GR with a minimally coupled scalar field with some potential.

In the absence of spatial curvature and other matter:

$$H^2 = \frac{\kappa^2}{3} \left(\frac{\dot{\phi}^2}{2} + V(\phi) \right)$$

$$\dot{H} = -\frac{\kappa^2}{2} \dot{\phi}^2$$

$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$$

where $\kappa^2 = 8\pi G$ ($\hbar = c = 1$).

Reduction to the first order equation

It can be reduced to the first order Hamilton-Jacobi-like equation for $H(\phi)$. From the equation for \dot{H} , $\frac{dH}{d\phi} = -\frac{\kappa^2}{2}\dot{\phi}$. Inserting this into the equation for H^2 , we get

$$\frac{2}{3\kappa^2} \left(\frac{dH}{d\phi} \right)^2 = H^2 - \frac{\kappa^2}{3} V(\phi)$$

Time dependence is determined using the relation

$$t = -\frac{\kappa^2}{2} \int \left(\frac{dH}{d\phi} \right)^{-1} d\phi$$

However, during oscillations of ϕ , $H(\phi)$ acquires non-analytic behaviour of the type $\text{const} + \mathcal{O}(|\phi - \phi_1|^{3/2})$ at the points where $\dot{\phi} = 0$, and then the correct matching with another solution is needed.

Inflationary slow-roll dynamics

Slow-roll occurs if: $|\ddot{\phi}| \ll H|\dot{\phi}|$, $\dot{\phi}^2 \ll V$, and then $|\dot{H}| \ll H^2$.

Necessary conditions: $|V'| \ll \kappa V$, $|V''| \ll \kappa^2 V$. Then

$$H^2 \approx \frac{\kappa^2 V}{3}, \quad \dot{\phi} \approx -\frac{V'}{3H}, \quad N \equiv \ln \frac{a_f}{a} \approx \kappa^2 \int_{\phi_f}^{\phi} \frac{V}{V'} d\phi$$

First obtained in A. A. Starobinsky, Sov. Astron. Lett. 4, 82 (1978) in the $V = \frac{m^2 \phi^2}{2}$ case and for a bouncing model.

Quantum generation of perturbations during inflation

Quantization with the adiabatic vacuum initial condition (in the tensor case, omitting the polarization tensor):

$$\hat{\phi} = (2\pi)^{-3/2} \int \left[\hat{a}_{\mathbf{k}} \phi_k(\eta) e^{-i\mathbf{k}\mathbf{r}} + \hat{a}_{\mathbf{k}}^\dagger \phi_k^* e^{i\mathbf{k}\mathbf{r}} \right] d^3k$$

where ϕ stands for ζ, g^a correspondingly and ϕ_k satisfies the equation

$$\frac{1}{f} (f \phi_k)'' + \left(k^2 - \frac{f''}{f} \right) \phi_k = 0, \quad \eta = \int \frac{dt}{a(t)}$$

For GW: $f = a$, for scalar perturbations in scalar field driven inflation in GR: $f = \frac{a\dot{\phi}}{H}$ where, in turn, the background scalar field satisfies the equation

$$\ddot{\phi} + 3H\dot{\phi} + \frac{dV}{d\phi} = 0$$

How the two basic hypothesis of the inflationary paradigm work.

I. Inflationary background: $t = \infty$ corresponds to $\eta = 0$ and $H(\eta) \equiv \frac{a'}{a^2}$ is bounded and slowly decreasing in this limit, so that $\frac{f''}{f} \sim \frac{2}{\eta^2}$. Then

$$\eta \rightarrow -0 : \quad \phi_k(\eta) \rightarrow \phi(k) = \text{const}, \quad P(k) = \frac{k^3 |\phi^2(k)|}{2\pi^2}$$

II. Adiabatic vacuum initial condition:

$$\eta \rightarrow -\infty : \quad \phi_k(\eta) = \frac{e^{-ik\eta}}{f\sqrt{2k}}$$

Combining both conditions:

$$\phi_k(\eta) \approx \frac{e^{-ik\eta}}{f\sqrt{2k}} \left(1 - \frac{i}{k\eta} \right), \quad a(\eta) \approx \frac{1}{H(\eta)|\eta|}$$

Spectral predictions of the one-field inflationary scenario in GR

Scalar (adiabatic) perturbations:

$$P_{\zeta}(k) = \frac{H_k^4}{4\pi^2 \dot{\phi}^2} = \frac{GH_k^4}{\pi |\dot{H}|_k} = \frac{128\pi G^3 V_k^3}{3V_k'^2}$$

where the index k means that the quantity is taken at the moment $t = t_k$ of the Hubble radius crossing during inflation for each spatial Fourier mode $k = a(t_k)H(t_k)$. Through this relation, the number of e-folds from the end of inflation back in time $N(t)$ transforms to $N(k) = \ln \frac{k_f}{k}$ where $k_f = a(t_f)H(t_f)$, t_f denotes the end of inflation.

The spectral slope

$$n_s(k) - 1 \equiv \frac{d \ln P_{\zeta}(k)}{d \ln k} = \frac{1}{\kappa^2} \left(2 \frac{V_k''}{V_k} - 3 \left(\frac{V_k'}{V_k} \right)^2 \right)$$

is small by modulus – confirmed by observations!

Tensor perturbations (A. A. Starobinsky, JETP Lett. 50, 844 (1979)):

$$P_g(k) = \frac{16GH_k^2}{\pi}; \quad n_g(k) \equiv \frac{d \ln P_g(k)}{d \ln k} = -\frac{1}{\kappa^2} \left(\frac{V'_k}{V_k} \right)^2$$

The consistency relation:

$$r(k) \equiv \frac{P_g}{P_\zeta} = \frac{16|\dot{H}_k|}{H_k^2} = 8|n_g(k)|$$

Tensor perturbations are always **suppressed** by at least the factor $\sim 8/N(k)$ compared to scalar ones. For the present Hubble scale, $N(k_H) = (50 - 60)$. Typically, $|n_g| \leq |n_s - 1|$, so $r \leq 8(1 - n_s) \sim 0.3$ – confirmed by observations!

Duration of inflation

Duration of inflation was finite inside our past light cone. In terms of e-folds, difference in its total duration in different points of space can be seen by the naked eye from a smoothed CMB temperature anisotropy map.

ΔN formalism: $\Delta\zeta(\mathbf{r}) = \Delta N_{tot}(\mathbf{r})$ where
 $N_{tot} = \ln \left(\frac{a(t_{fin})}{a(t_{in})} \right) = N_{tot}(\mathbf{r})$ (AS, 1982,1985).

For $\ell \lesssim 50$, neglecting the Silk and Doppler effects, as well as the ISW effect due the presence of dark energy,

$$\frac{\Delta T(\theta, \phi)}{T_\gamma} = -\frac{1}{5} \Delta\zeta(r_{LSS}, \theta, \phi) = -\frac{1}{5} \Delta N_{tot}(r_{LSS}, \theta, \phi)$$

For $\frac{\Delta T}{T} \sim 10^{-5}$, $\Delta N \sim 5 \times 10^{-5}$, and for $H \sim 10^{14}$ GeV,
 $\Delta t \sim 5t_{Pl}$!

Inverse reconstruction of inflationary models in GR

In the slow-roll approximation:

$$\frac{V^3}{V'^2} = C P_\zeta(k(t(\phi))), \quad C = \frac{12\pi^2}{\kappa^6}$$

Changing variables for ϕ to $N(\phi)$ and integrating, we get:

$$\frac{1}{V(N)} = -\frac{\kappa^4}{12\pi^2} \int \frac{dN}{P_\zeta(N)}$$

$$\kappa\phi = \int dN \sqrt{\frac{d \ln V}{dN}}$$

Here, $N \gg 1$ stands both for $\ln(k_f/k)$ at the present time and the number of e-folds back in time from the end of inflation. First derived in H. M. Hodges and G. R. Blumenthal, Phys. Rev. D 42, 3329 (1990).

The two-parameter family of **isospectral** slow-roll inflationary models, but the second parameter shifts the field ϕ only.

Minimal "scale-free" reconstruction

Minimal inflationary model reconstruction avoiding introduction of any new physical scale **both** during and after inflation and producing the best fit to the Planck data.

Assumption: the numerical coincidence between $2/N_H \sim 0.04$ and $1 - n_s$ is not accidental but happens for all $1 \ll N \lesssim 60$: $P_\zeta = P_0 N^2$. Then:

$$V = V_0 \frac{N}{N + N_0} = V_0 \tanh^2 \frac{\kappa \phi}{2\sqrt{N_0}}$$

$$r = \frac{8N_0}{N(N + N_0)}$$

$r \sim 0.003$ for $N_0 \sim 1$. From the upper limit on r :

$$N_0 < \frac{0.07N^2}{8 - 0.07N}$$

$N_0 < 57$ for $N = 57$.

Another example: $P_\zeta = P_0 N^{3/2}$.

$$V(\phi) = V_0 \frac{\phi^2 + 2\phi\phi_0}{(\phi + \phi_0)^2}$$

Not bounded from below (of course, in the region where the slow-roll approximation is not valid anymore). Crosses zero linearly.

More generally, the two "aesthetic" assumptions – "no-scale" scalar power spectrum and $V \propto \phi^{2n}$, $n = 1, 2, \dots$ at the minimum of the potential – lead to

$P_\zeta = P_0 N^{n+1}$, $n_s - 1 = -\frac{n+1}{N}$ unambiguously. From this, only $n = 1$ is permitted by observations. Still an additional parameter appears due to tensor power spectrum – no preferred one-parameter model (if the $V(\phi) \propto \phi^2$ model is excluded).

Inflation in $f(R)$ gravity

The simplest model of modified gravity (= geometrical dark energy) considered as a phenomenological macroscopic theory in the fully non-linear regime and non-perturbative regime.

$$S = \frac{1}{16\pi G} \int f(R) \sqrt{-g} d^4x + S_m$$

$$f(R) = R + F(R), \quad R \equiv R^\mu{}_\mu$$

Here $f''(R)$ is not identically zero. Usual matter described by the action S_m is minimally coupled to gravity.

Vacuum one-loop corrections depending on R only (not on its derivatives) are assumed to be included into $f(R)$. The normalization point: at laboratory values of R where the scalaron mass (see below) $m_s \approx \text{const.}$

Metric variation is assumed everywhere. Palatini variation leads to a different theory with a different number of degrees of freedom.

Field equations

$$\frac{1}{8\pi G} \left(R^\nu_\mu - \frac{1}{2} \delta^\nu_\mu R \right) = - \left(T^\nu_{\mu(vis)} + T^\nu_{\mu(DM)} + T^\nu_{\mu(DE)} \right) ,$$

where $G = G_0 = \text{const}$ is the Newton gravitational constant measured in laboratory and the effective energy-momentum tensor of DE is

$$8\pi G T^\nu_{\mu(DE)} = F'(R) R^\nu_\mu - \frac{1}{2} F(R) \delta^\nu_\mu + (\nabla_\mu \nabla^\nu - \delta^\nu_\mu \nabla_\gamma \nabla^\gamma) F'(R) .$$

Because of the need to describe DE, de Sitter solutions in the absence of matter are of special interest. They are given by the roots $R = R_{dS}$ of the algebraic equation

$$Rf'(R) = 2f(R) .$$

The special role of $f(R) \propto R^2$ gravity: admits de Sitter solutions with **any** curvature.

Degrees of freedom

I. In quantum language: particle content.

1. **Graviton** – spin 2, massless, transverse traceless.

2. **Scalaron** – spin 0, massive, mass - R -dependent:

$$m_s^2(R) = \frac{1}{3f''(R)} \text{ in the WKB-regime.}$$

II. Equivalently, in classical language: number of free functions of spatial coordinates at an initial Cauchy hypersurface.

Six, instead of four for GR – two additional functions describe massive scalar waves.

Thus, $f(R)$ gravity is a **non-perturbative** generalization of GR.

It is equivalent to scalar-tensor gravity with $\omega_{BD} = 0$ if

$$f'''(R) \neq 0.$$

Transformation to the Einstein frame and back

In the Einstein frame, free particles of usual matter do not follow geodesics and atomic clocks do not measure proper time.

From the Jordan (physical) frame to the Einstein one:

$$g_{\mu\nu}^E = f' g_{\mu\nu}^J, \quad \kappa\phi = \sqrt{\frac{3}{2}} \ln f', \quad V(\phi) = \frac{f'R - f}{2\kappa^2 f'^2}$$

where $\kappa^2 = 8\pi G$.

Inverse transformation:

$$R = \left(\sqrt{6}\kappa \frac{dV(\phi)}{d\phi} + 4\kappa^2 V(\phi) \right) \exp \left(\sqrt{\frac{2}{3}} \kappa\phi \right)$$
$$f(R) = \left(\sqrt{6}\kappa \frac{dV(\phi)}{d\phi} + 2\kappa^2 V(\phi) \right) \exp \left(2\sqrt{\frac{2}{3}} \kappa\phi \right)$$

$V(\phi)$ should be at least C^1 .

Reduction to the first order equation

In the absence of spatial curvature and $\rho_m = 0$, it is always possible to reduce these equations to a first order one using either the transformation to the Einstein frame and the Hamilton-Jacobi-like equation for a minimally coupled scalar field in a spatially flat FLRW metric, or by directly transforming the 0-0 equation to the equation for $R(H)$:

$$\frac{dR}{dH} = \frac{(R - 6H^2)f'(R) - f(R)}{H(R - 12H^2)f''(R)}$$

Analogues of large-field (chaotic) inflation: $F(R) \approx R^2 A(R)$ for $R \rightarrow \infty$ with $A(R)$ being a slowly varying function of R , namely

$$|A'(R)| \ll \frac{A(R)}{R} \ , \ |A''(R)| \ll \frac{A(R)}{R^2} \ .$$

Analogues of small-field (new) inflation, $R \approx R_1$:

$$F'(R_1) = \frac{2F(R_1)}{R_1} \ , \ F''(R_1) \approx \frac{2F(R_1)}{R_1^2} \ .$$

Thus, all inflationary models in $f(R)$ gravity are close to the simplest one over some range of R .

Perturbation spectra in slow-roll $f(R)$ inflationary models

Let $f(R) = R^2 A(R)$. In the slow-roll approximation $|\ddot{R}| \ll H|\dot{R}|$:

$$P_\zeta(k) = \frac{\kappa^2 A_k}{64\pi^2 A_k'^2 R_k^2}, \quad P_g(k) = \frac{\kappa^2}{12A_k\pi^2}$$

$$N(k) = -\frac{3}{2} \int_{R_f}^{R_k} dR \frac{A}{A'R^2}$$

where the index k means that the quantity is taken at the moment $t = t_k$ of the Hubble radius crossing during inflation for each spatial Fourier mode $k = a(t_k)H(t_k)$.

Smooth reconstruction of inflation in $f(R)$ gravity

$$f(R) = R^2 A(R)$$

$$A = \text{const} - \frac{\kappa^2}{96\pi^2} \int \frac{dN}{P_\zeta(N)}$$

$$\ln R = \text{const} + \int dN \sqrt{-\frac{2}{3} \frac{d \ln A}{dN}}$$

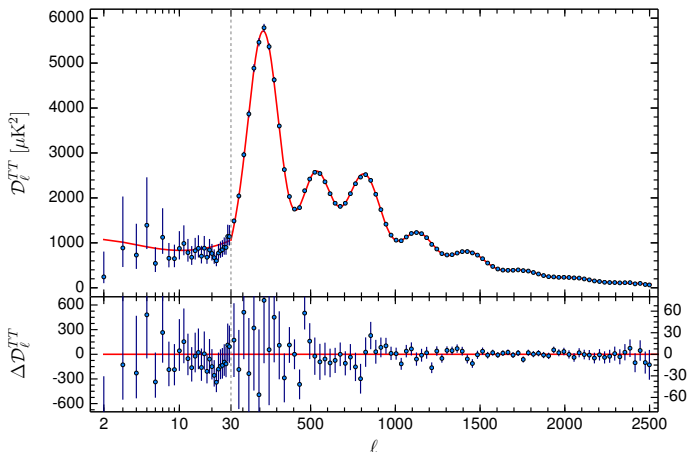
Here, the additional assumptions that $P_\zeta \propto N^\beta$ and that the resulting $f(R)$ can be analytically continued to the region of small R without introducing a new scale, and it has the linear (Einstein) behaviour there, leads to $\beta = 2$ and the $R + R^2$ inflationary model with $r = \frac{12}{N^2} = 3(n_s - 1)^2$ unambiguously.

Perspectives of future discoveries

- ▶ Primordial gravitational waves from inflation: r .
 $r \lesssim 8(1 - n_s) \approx 0.3$ (confirmed!) but may be much less.
- ▶ Local non-smooth features in the scalar power spectrum at cosmological scales (?).
- ▶ Local enhancement of the power spectrum at small scales leading to a significant amount of primordial black holes (?).

Small features in the power spectrum

- 1) A $\sim 10\%$ depression for $20 \lesssim \ell \lesssim 40$.
- 2) An upward wiggle at $\ell \approx 40$ (the Archeops feature) and a downward one at $\ell \approx 22$.



Local features in an inflaton potential

Some new physics beyond one slow-rolling inflaton may show itself through these features.

The simplest models with two additional parameters which can describe such behaviour (to some extent) are based on the exactly soluble model considered in [A. A. Starobinsky, JETP Lett. **55**, 489 \(1992\)](#): an inflaton potential with a sudden change of its first derivative.

For comparison of some elaborated class of such models with the CMB TT data, see

[D. K. Hazra, A. Shafieloo, G. F. Smoot and A. A. Starobinsky, JCAP **1408**, 048 \(2014\)](#).

Recent comparison of the model

$$V(\phi) = \theta(\phi_0 - \phi) V_1 (1 - \exp(-\alpha\kappa\phi)) + \\ \theta(\phi - \phi_0) V_2 (1 - \exp(-\alpha\kappa(\phi - \phi_1)))$$

where

$$V_1 (1 - \exp(-\alpha\kappa\phi_0)) = V_2 (1 - \exp(-\alpha\kappa(\phi_0 - \phi_1)))$$

with the Planck2015 data on both CMB temperature anisotropy and E-mode polarization (D. K. Hazra, A. Shafieloo, G. F. Smoot and A. A. Starobinsky, JCAP **1606**, 007 (2016)) provides ~ 12 improvement in χ^2 fit to the data compared to best-fit models with smooth inflaton potentials, partly due to the better description of wiggles at both $\ell \approx 40$ and $\ell \approx 22$. Possible confusion effects, in particular due to an extended ionization history, have been recently analyzed in D. K. Hazra, D. Paoletti, M. Ballardini, F. Finelli, A. Shafieloo, G. F. Smoot and A. A. Starobinsky, JCAP **1802**, 017 (2018); arXiv:1710.01205.

Conclusions

- ▶ The typical inflationary predictions that $|n_s - 1|$ is small and of the order of N_H^{-1} , and that r does not exceed $\sim 8(1 - n_s)$ are confirmed. Typical consequences following without assuming additional small parameters: $H_{55} \sim 10^{14}$ GeV, $m_{infl} \sim 10^{13}$ GeV.
- ▶ Though the Einstein gravity plus a minimally coupled inflaton remains sufficient for description of inflation with existing observational data, modified (in particular, scalar-tensor or $f(R)$) gravity can do it as well.
- ▶ From the scalar power spectrum $P_\zeta(k)$, it is possible to reconstruct an inflationary model both in the Einstein and $f(R)$ gravity up to one arbitrary physical constant of integration.

- ▶ In the Einstein gravity, the simplest inflationary models permitted by observational data are two-parametric, no preferred quantitative prediction for r , apart from its parametric dependence on $n_s - 1$, namely, $\sim (n_s - 1)^2$ or larger.
- ▶ In the $f(R)$ gravity, the simplest model is one-parametric and has the preferred value $r = \frac{12}{N^2} = 3(n_s - 1)^2$.
- ▶ Thus, it has sense to search for primordial GW from inflation at the level $r > 10^{-3}$!
- ▶ Investigation of small local features in the primordial power spectrum (bispectrum, etc.) of scalar perturbations can provide "tomography" of inflation and may lead to discovery of other particles or quasiparticles more massive than inflaton.